

Remote state preparation based on optical hybrid entanglement

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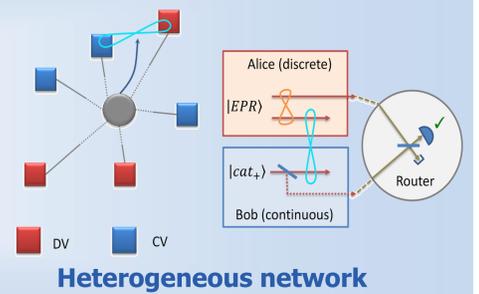


Introduction

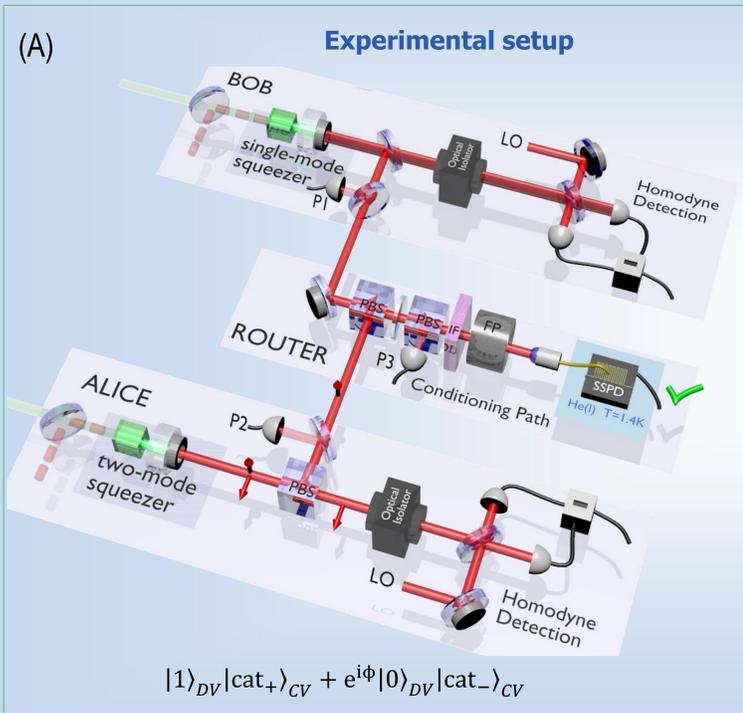
The wave-particle duality of light has led to two different encodings for optical quantum information processing. Several approaches have emerged based either on particle-like discrete-variable states, e.g. finite-dimensional quantum systems, or on wave-like continuous-variable states, e.g. infinite-dimensional systems.

Aiming at connecting the two worlds in order to build quantum heterogeneous networks [1], where a discrete qubit can be converted in a continuous encoding, we demonstrate the first measurement-induced generation of hybrid entanglement between optical quantum bits of these different types [2], located at distant places and connected by a lossy channel.

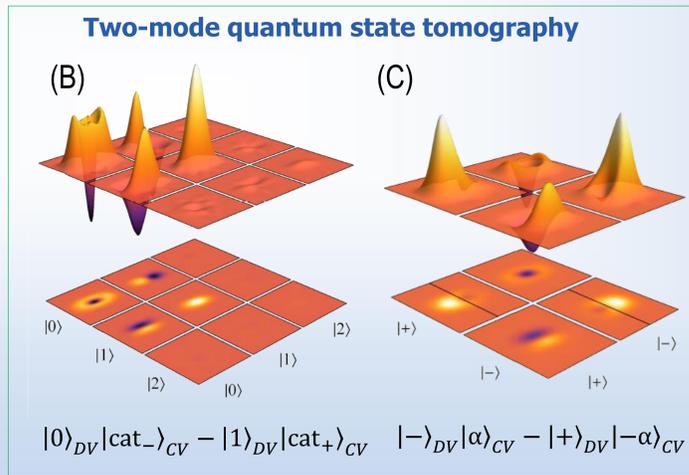
Using this hybrid resource, we can engineer remotely any coherent state superposition or qu-mode by conditioning on the results of quadrature measurements performed on the discrete part of this hybrid state



Measurement-induced generation of hybrid entanglement

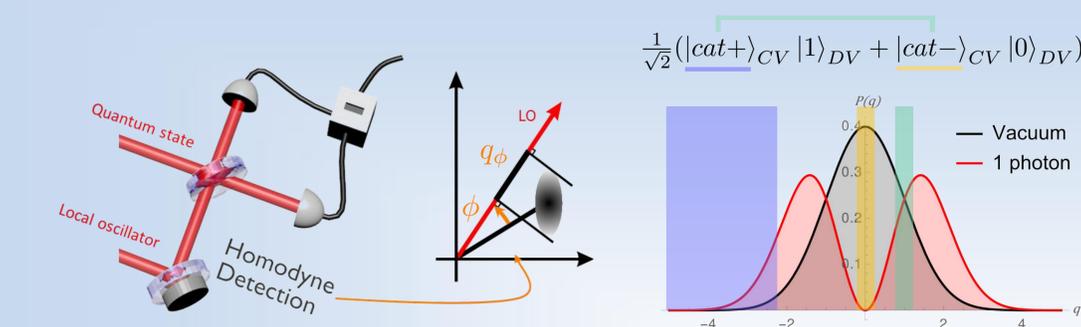


(A) Alice and Bob locally generate the required resources by using continuous-wave optical parametric oscillators (OPOs) operated below threshold [3][4]. A two-mode squeezer and a single-mode squeezer are used respectively on Alice's and Bob's node. A small fraction of Bob's squeezed vacuum is tapped (3%) and mixed at a central station to the idler beam generated by Alice. The resulting beam is then frequency filtered and detected by a superconducting single photon detector (SSSPD). Given a detection event, which heralds the generation of entanglement between an even and an odd cat state and the presence or absence of a single photon, the entangled state is characterized by two high-efficiency homodyne detections in the optimal temporal modes [5].



(B) The beam splitter ratio in the central station is adjusted to generate a maximally entangled state. The figure gives Wigner functions associated with the reduced density matrices $\langle k|\hat{\rho}|l\rangle$ with $k, l \in \{0, 1, 2\}$ after correction for detection losses ($\eta = 85\%$). For components with $k \neq l$, the plots give the real part of $\langle k|\hat{\rho}|l\rangle$ (back corner) and the imaginary part of $\langle l|\hat{\rho}|k\rangle$ (front corner). **(C)** Wigner functions associated with the reduced density matrices $\langle k|\hat{\rho}|l\rangle$ with $k, l \in \{+, -\}$, where $|\pm\rangle = (|0\rangle \pm |1\rangle)/\sqrt{2}$. A fidelity $77 \pm 3\%$ is obtained with the target state with $|\alpha| = 0.9$.

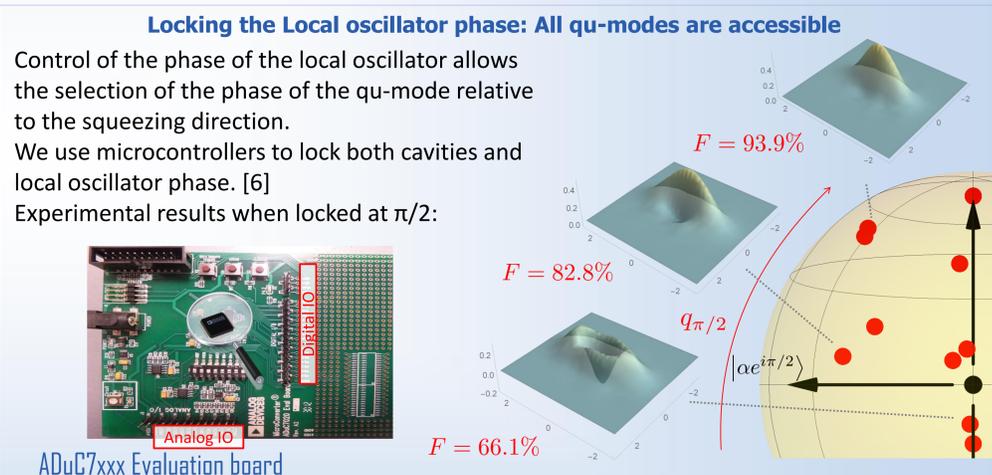
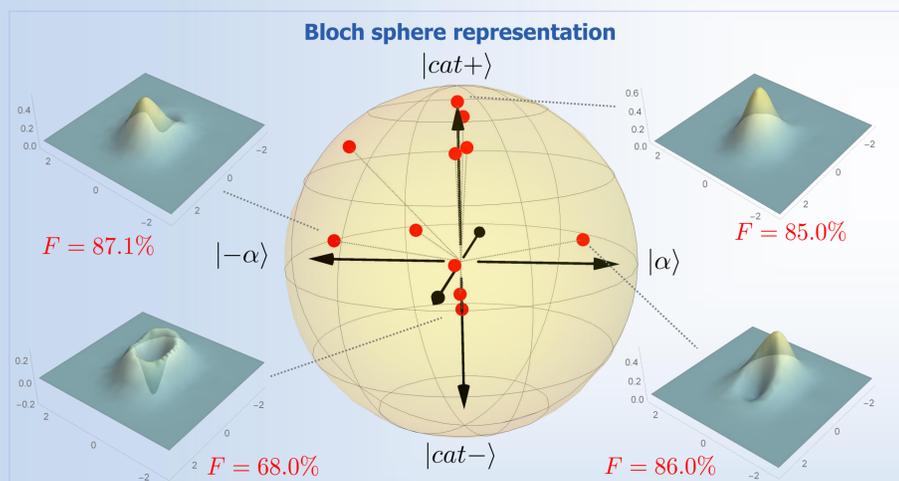
Remote state preparation: Generation of arbitrary continuous variable qu-modes



Conditioning on the result of the homodyne measurement on the DV side, it is possible to choose the state created on the CV side.

Any qu-mode is accessible by choosing ϕ and q_ϕ .

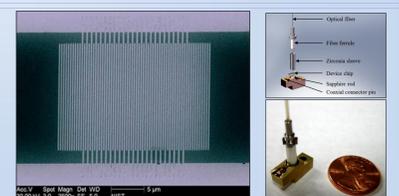
$$\langle q_\phi | \Psi \rangle_{DV} = \frac{1}{\sqrt{N_q}} (|cat_-\rangle_{CV} + q_\phi e^{i\phi} |cat_+\rangle_{CV})$$



Superconducting Nanowire Single-photon detector (SNSPD)



To herald our state generation, we use SNSPDs based on WSi from a collaboration with the National Institute of Standards and Technology (NIST) and the Jet Propulsion Laboratory (JPL) [7]. Those detectors reach 93% of detection efficiency at 1064 nm [8] instead of 1% for an avalanche photodiode, with low dark counts ($<10\text{Hz}$). Using them we characterise our single photon source as a 1 MHz generation rate and above 90% heralding efficiency.



References:

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