

Quantum enhanced classical imaging and metrology

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University-company collaboration



Palacky University- Optics department



Meopta-Optika,
2500 employees
150 people R&D

Quantum state tomography formalism

- The goal is to estimate the quantum state from the measurement data obtained from the ensemble of N identical copies the quantum system
- Measurements are described by a set of positive operators Π_j (POVM operators)
- Due to finite resources, quantum state can be only statistically estimated from probabilities $p_j = \text{Tr } \rho \Pi_j$
- The probabilities are measured by the outcome frequencies f_j of the particular measurements $f_j = \frac{n_j}{N}$

Fisher information and estimation

How good is my estimator of ς in the case of $\varrho(\varsigma)$? Any unbiased estimator must follow Cramer-Rao bound:

$$(\Delta\hat{\varsigma})^2 \geq \frac{1}{\mathcal{F}}$$

Classical Fisher information: $\mathcal{F} = \int_{-\infty}^{\infty} \varrho_{\varsigma}(x) \left(\frac{\partial \log \varrho_{\varsigma}(x)}{\partial \varsigma} \right)^2 dx$.

Quantum Fisher information: $\mathcal{F}_{\mathcal{Q}} = \text{Tr}[\varrho_{\varsigma} L_{\varsigma}^2]$

symmetric logarithmic derivative L_{ς} is the selfadjoint operator satisfying $\frac{1}{2}(L_{\varsigma}\varrho_{\varsigma} + \varrho_{\varsigma}L_{\varsigma}) = \partial\varrho_{\varsigma}/\partial\varsigma$

Fisher information notes

Quantum Fisher information is an upper bound for a classical Fisher information,

$$\mathcal{F} \geq \mathcal{F}_{\mathcal{Q}}.$$

Quantum fisher information is independent of a measurement process. To test optimality of particular measurement, Fisher information for the measurement has to be computed:

$$\mathcal{F}_{\mathcal{M}} \geq \mathcal{F}_{\mathcal{Q}}.$$

Correspondence classical and quantum description

There is a tight correspondence between beam optics and quantum mechanics.

- coherent waves (beam modes) → pure states

$$U(x) = \langle x | \psi \rangle$$

- partially coherent fields → mixed states

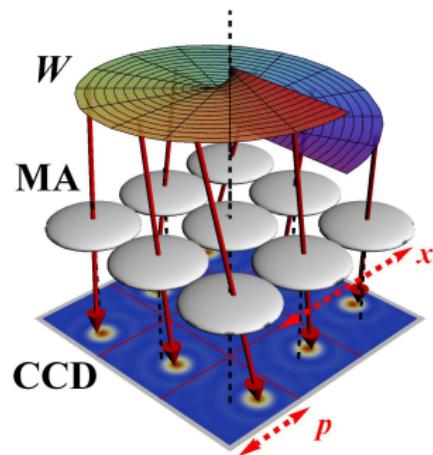
$$G(x, x') = \langle x | \rho | x' \rangle$$

Wavefront detector tomography

Measurement of optical beams spatial coherence parametrized by a coherence matrix ϱ

$$I(\Delta x_i, \Delta p_j) = \text{Tr}(\varrho |\Pi_{ij}\rangle \langle \Pi_{ij}|)$$

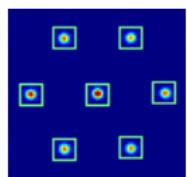
$$(\Pi_{ij})_{mn} = \psi_{n,i}(\Delta p_j) \psi_{m,i}^*(\Delta p_j)$$



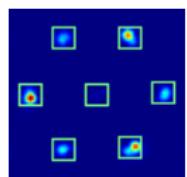
Vortex beam reconstruction

Vortex reconstruction space

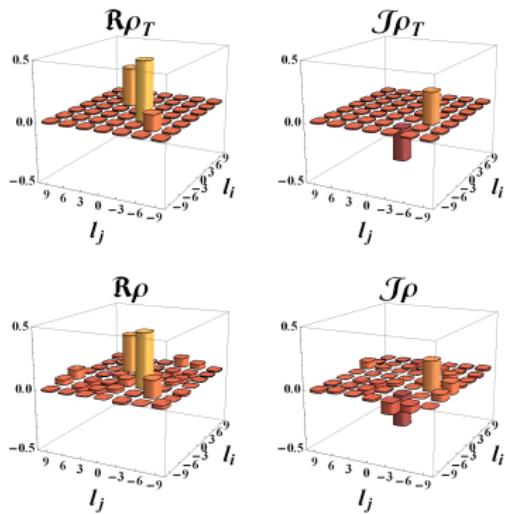
$$V_\ell = \langle r, \varphi | V_\ell \rangle \propto e^{i\ell\varphi}$$
$$\ell \in \{-9, -6, -3, 0, +3, +6, +9\}$$



a

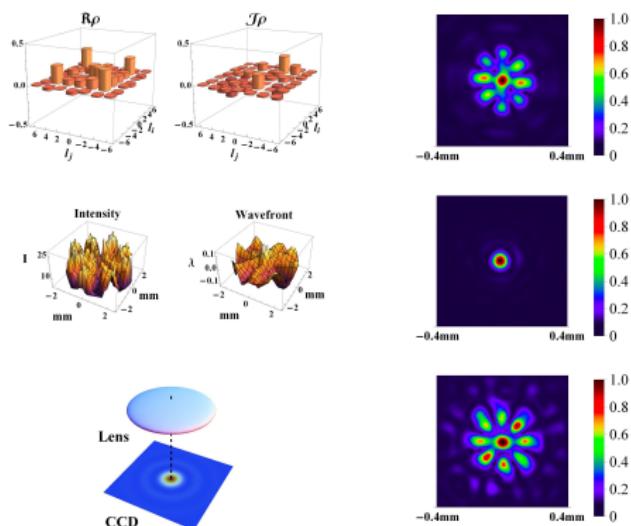


b



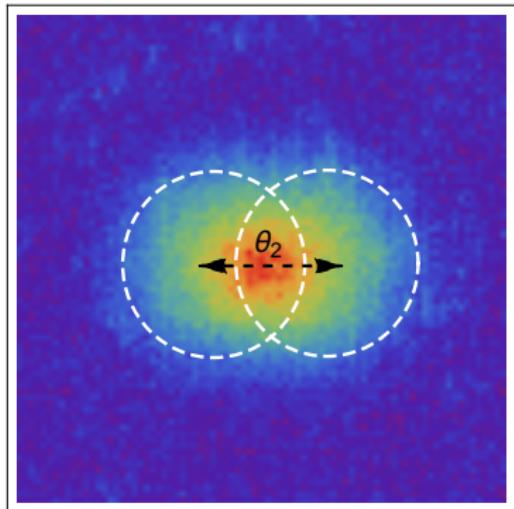
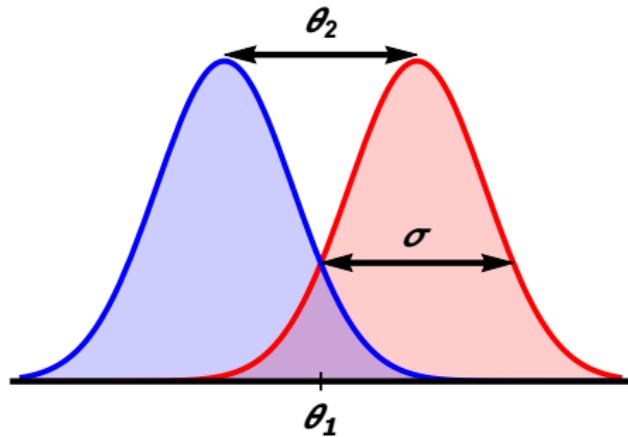
Partially coherent light intensity propagation

The propagation of transverse intensity distribution requires the explicit form of mutual coherence function at the input plane



Incoherent image of two-points

$$\text{PSF: } I(x) = |\langle x | \psi \rangle|^2 = |\psi(x)|^2$$



$$|\psi_{\pm}\rangle = \exp(\pm iP\mathfrak{s}/2)|\psi\rangle, \varrho_{\mathfrak{s}} = \frac{1}{2}(|\psi_+\rangle\langle\psi_+| + |\psi_-\rangle\langle\psi_-|)$$

Recent experimental work

Taking resolution to the limit: dispelling Rayleigh curse

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Optimal measurement

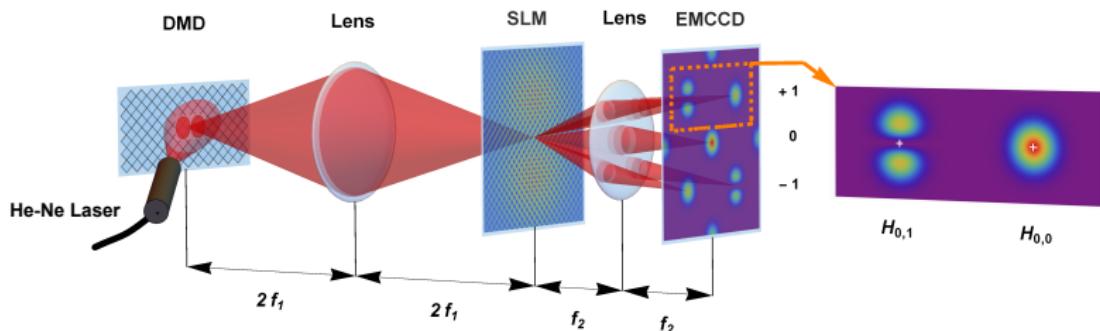
$$|\psi_{\pm}\rangle = \exp(\pm iP\mathfrak{s}/2)|\psi\rangle, \langle\psi_-|\psi_+\rangle \neq 0$$

$$\begin{aligned} |\psi_{sm}\rangle &= C_{sm}(|\psi_+\rangle + |\psi_-\rangle) \simeq |\psi\rangle, \\ |\psi_a\rangle &= C_a(|\psi_+\rangle - |\psi_-\rangle) \simeq \frac{P|\psi\rangle}{\sqrt{\langle\psi|P^2|\psi\rangle}}, \end{aligned}$$

Once PSF is inversion symmetric, those modes are orthogonal and Quantum Fisher information is:

$$\mathcal{F}_{\mathcal{Q}} = 2 \left[\frac{1}{p_a} \langle\psi_a| \frac{\partial\varrho_{\mathfrak{s}}}{\partial\mathfrak{s}} |\psi_a\rangle + \frac{1}{p_{sm}} \langle\psi_{sm}| \frac{\partial\varrho_{\mathfrak{s}}}{\partial\mathfrak{s}} |\psi_{sm}\rangle \right] \simeq \langle\psi|P^2|\psi\rangle,$$

Experimental realization of mode projection



In the direction of the hologram reference wave, observed intensity is:

$$\left(\left| \int_{-\infty}^{\infty} dx \phi_q^* \phi_0(x + \frac{\theta_2}{2}) \right|^2 + \left| \int_{-\infty}^{\infty} dx \phi_q^* \phi_0(x - \frac{\theta_2}{2}) \right|^2 \right)$$

Benefits of quantum description of measurement and imaging

- Recasting classical measurement scenario should provide a new point of view about a problem (Shack-Hartmann example).
- Proper treatment of Quantum Fisher Information provides real boundaries to measurement process and should lead to improvement in a measurement scheme (the image of two incoherent sources example).