CONDITIONAL PREPARATION OF LOW-NOISE BRIGHT TWIN BEAMS

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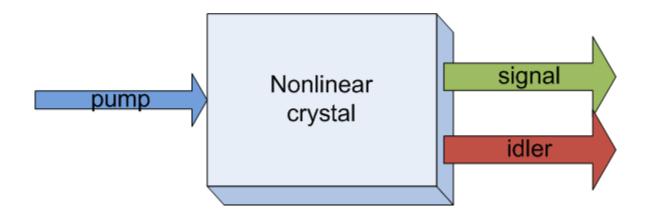
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Outline

- Twin-beam states and their use in CV quantum communication
- Multimode structure
- Feed-forward technique
- Increasing correlations between the beams
- Reducing photon-number fluctuations within the beams
- Summary

Twin-beam states (TWB)

Output of an OPA with a vacuum input:



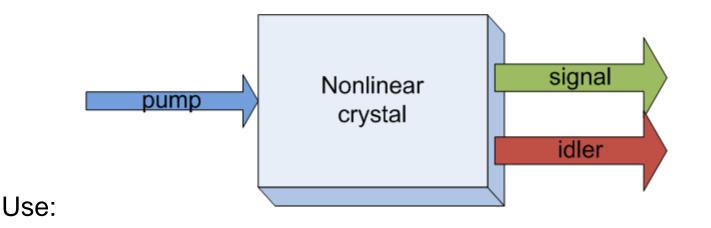
Pump photon is splitted to signal and idler photons.

They satisfy momentum and energy conservation.

The process can be described by the quadratic Hamiltonian.

Twin-beam states (TWB)

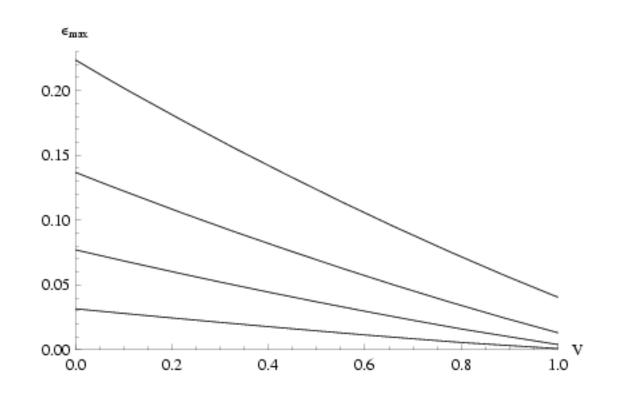
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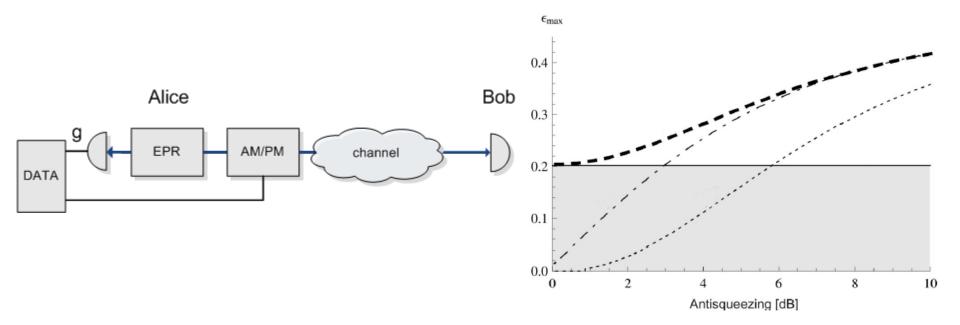
- Quantum metrology [G. Brida et al., J. Opt. Soc. Am. B 23, 2185 (2006)]
- Quantum imaging [G. Brida et al., Nature Photonics 4, 227 230 (2010)]
- Quantum key distribution [L. S. Madsen, VU et al., Nature Communications 3, 1083 (2012)]

- Squeezed states can partially substitute the inefficient classical error-correction in CV QKD [VU, R. Filip, NJP 13, 113007, 2011]
- Modulated TWB states can improve CV QKD [L. S. Madsen, VU et al., Nature Communications 3, 1083 (2012)]
- Multimode TWB can be used for CV QKD [VU, L. Ruppert, R. Filip, PRA 90, 062326, 2014]
- Optimally modulated squeezed states cancel information leakage in CV quantum channels [*C. S. Jacobsen, VU et al., arXiv:1408.4566*]

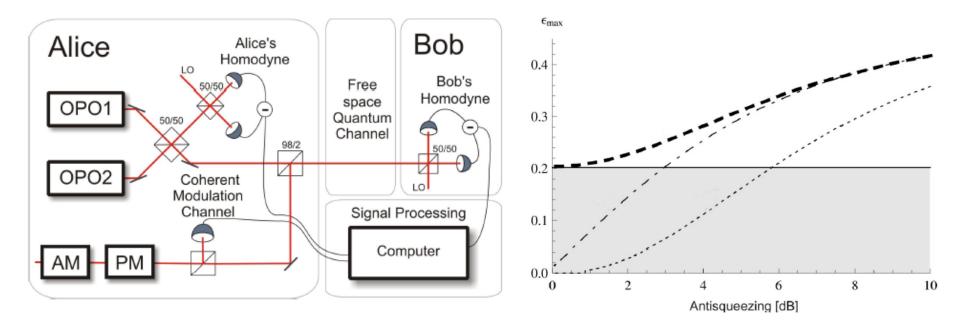
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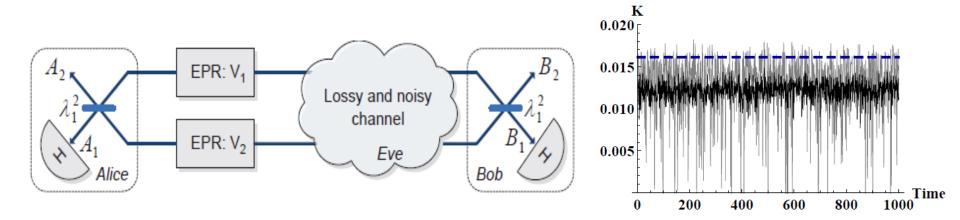
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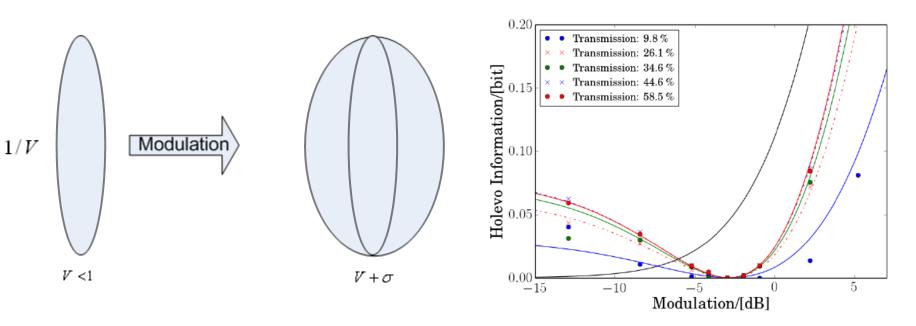


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It was about quadrature squeezing, now let's look at photon-number / intensity measurements.

TWB: photon-number description

Photon numbers $(a^{+}a)$ can be measured in each beam.

Representation in the Fock (number) basis:

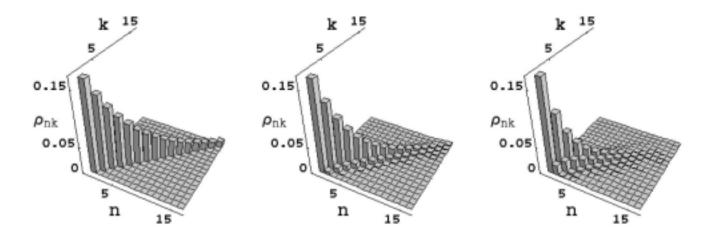
$$|x\rangle\rangle = \sqrt{(1-x^2)}\sum_n x^n \, |n,n\rangle\rangle \ , \ |n,n\rangle\rangle \equiv |n\rangle_s \otimes |n\rangle_i$$

 $x \in \mathbb{C}$ and $0 \leq |x| \leq 1$

 $x = \sqrt{N/(N+1)}$, where N – mean photon number

TWB: photon-number description

Photon-number statistics in each beam is thermal:

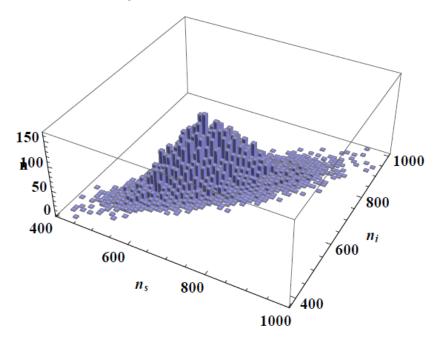


Correlation is (almost) perfect, thus **sub-shot-noise** photonnumber difference correlation is observed.

Bright twin-beams are heavily multimode (spatial, frequency etc)

The state becomes the mixture of multiple twin-beam states.

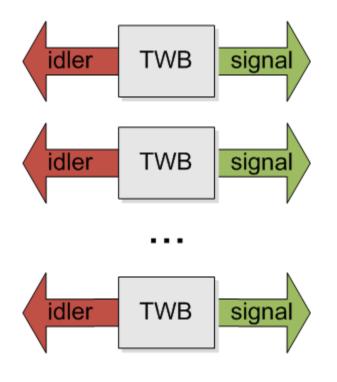
Statistics becomes "bell-shaped":



100 modes with N=10 photon each

Twin-beams can be multimode (spatial, frequency etc)

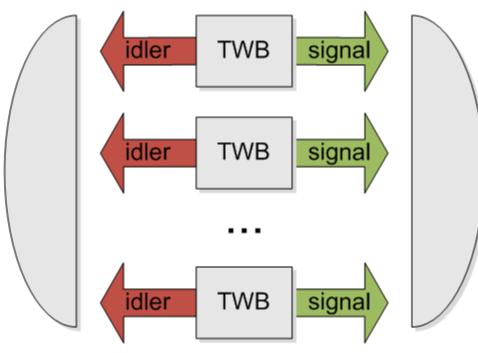
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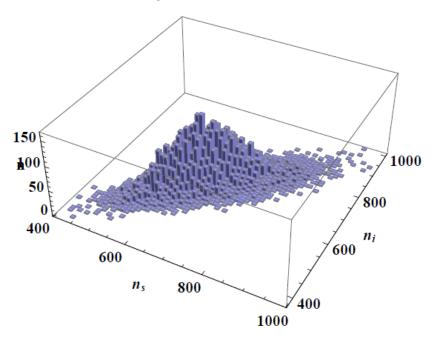
Mode-nondiscriminating photon-counting:



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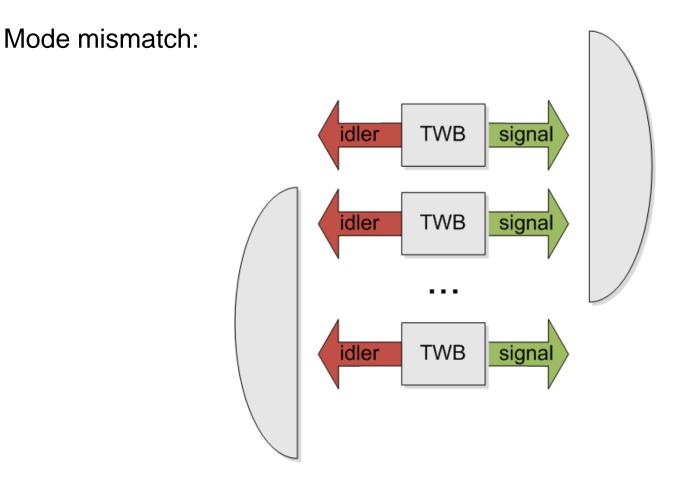
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Correlations are degraded due to losses, inefficient detection and mode mismatch. Nonclassicality can be lost.

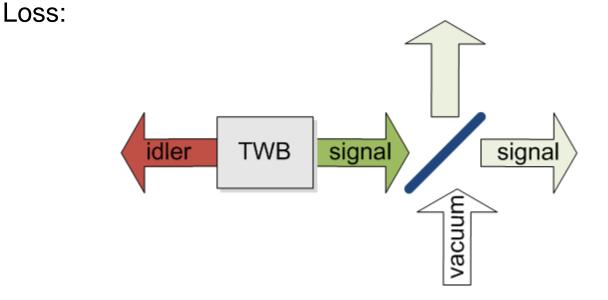
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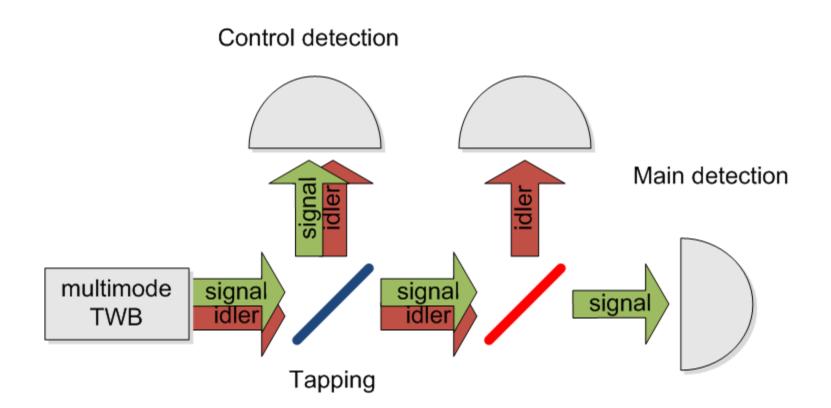
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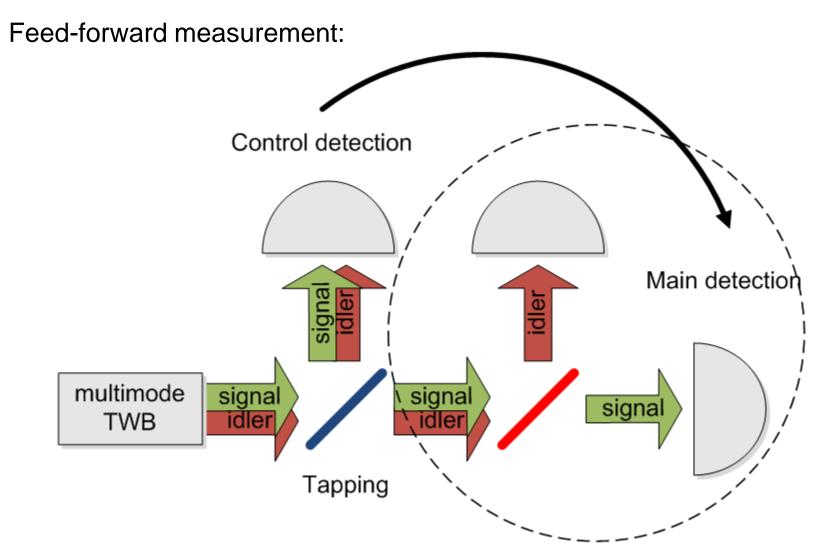


Restoring nonclassical correlations

Feed-forward measurement:



Restoring nonclassical correlations



When measurement outcome on control detector satisfies a condition (defined by position and width), the main detection is performed.

Restoring nonclassical correlations

Correlation can be characterized by noise reduction factor (NRF): $NRF \equiv Var(N_i - N_s)/\langle N_i + N_s \rangle$

In case of m matched and k mismatched modes:

$$NRF_{meas} = 1 - \frac{m}{m+k}\eta + \frac{k}{m+k}\eta N_{mode}.$$

Photon-number sum fluctuations can be characterized by normalized variance of photon-number sum:

 $F \equiv Var(N_i + N_s) / \langle N_i + N_s \rangle$

which reads $F = 2\eta N_{mode} + 1$.

Multimode TWB: Numerical model

Conditional preparation was modeled as generation of multimode TWB states, coupling to vacuum modes, mode-nondiscriminating photon-counting measurements and conditioning.

Generated states: $p(n) = (1 - x^2)x^{2n}$

After coupling to vacuum:

$$p(k) = \frac{n!}{k!(n-k)!}T^{n-k}(1-T)^k$$

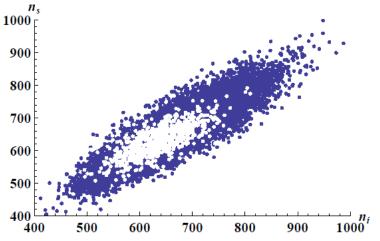
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Photon-number distributions before and after optimal conditioning

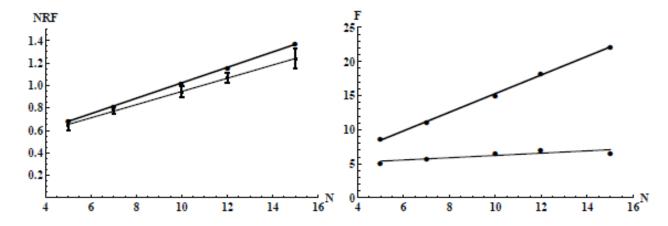
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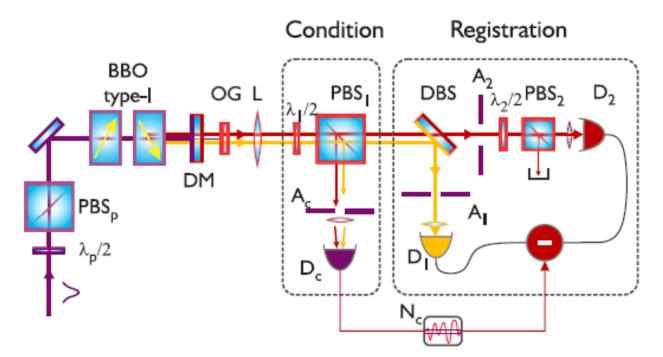
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NRF and F-factors before and after optimal conditioning

Multimode TWB: NRF experiment

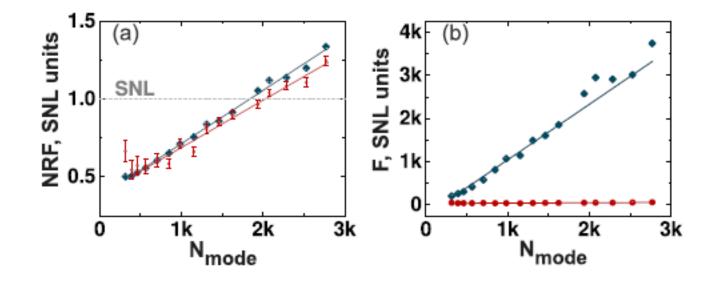
Correlation can be improved by feed-forward measurement:



Sketch of the experiment at MPI in Erlangen

Multimode TWB: NRF experiment

Experimental results:



Optimal conditioning leads to reduction of NRF and F factors.

Singe-mode case

Single-mode state:
$$|\Psi\rangle = \frac{1}{\sqrt{1 + N_{mode}}} \sum_{n=0}^{\infty} \lambda^{\frac{n}{2}} |n\rangle_{S} |n\rangle_{I}$$

with $\lambda = \frac{N_{mode}}{N_{mode}+1}$

After perfect subtraction of N photons:

$$(a_S a_I)^N |\Psi\rangle = \frac{1}{\sqrt{1 + N_{mode}}} \sum_{n=0}^{\infty} \frac{(n+N)!}{N!} \lambda^{\frac{n+N}{2}} |n\rangle_S |n\rangle_I$$

After normalization:

$$|\Psi'\rangle = \frac{\sum_{n=0}^{\infty} \frac{(n+N)!}{n!N!} \lambda^{\frac{n}{2}} |n\rangle_S |n\rangle_I}{\sqrt{{}_2F_1 \left[1+N, 1+N, 1, \lambda\right]}}$$

With increasing N the maximum moves away from the origin, brightness increases.

Statistics characterization

To characterize the statistics after conditioning, we suggest using mean-deviation-ratio (MDR):

$$\mathrm{MDR} = \frac{\langle n \rangle}{\sqrt{\langle (\Delta n)^2 \rangle}} \approx \begin{cases} \sqrt{2\lambda N - 1} & \lambda \ll 1 \text{ and } \operatorname{large} N, \\ \sqrt{2\lambda N + 1} & \lambda \to 1 \text{ and } \operatorname{large} N. \end{cases}$$

In the multimode regime MDR can grow with increasing number of modes. By conditional preparation we increase it without changing the mode structure.

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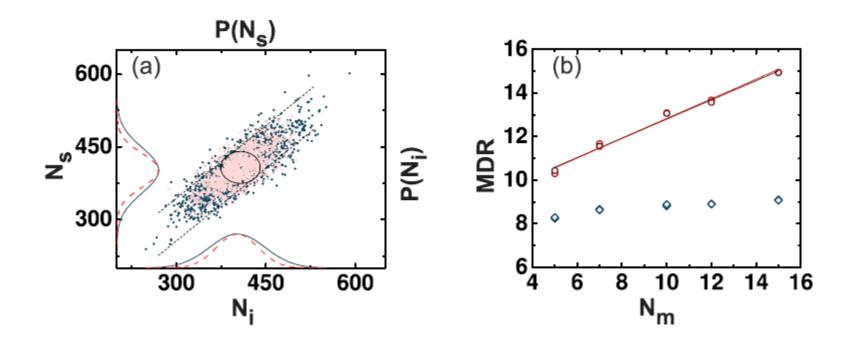
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Why not
$$g^{(2)} \equiv (\langle n^2 \rangle - \langle n \rangle) / \langle n \rangle^2$$
?

It can be decreased by the number of modes $g^{(2)}(0) = 1 + rac{1}{M}$

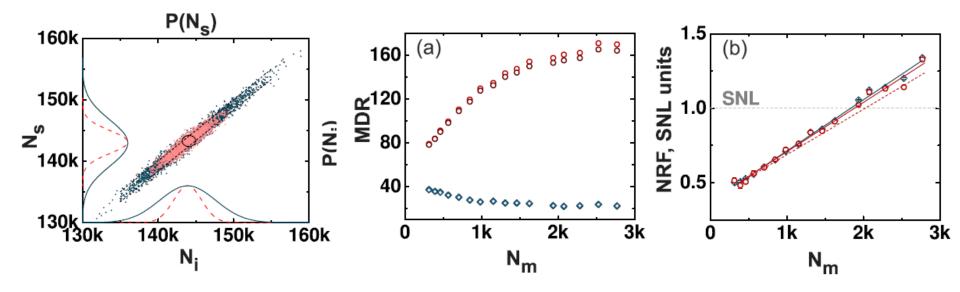
Why not Fano? It can be reduced by attenuation, it is weakly changed for intense beams.

Suppressing the noise: numerics



Distributions and MDR before and after conditioning

Suppressing the noise: experiment



Distributions, MDR and NRF before and after conditioning

Summary

• We suggest the method of conditional preparation of special low-noise twin-beam states;

• Numerical model and experimental test confirm possibility to increase correlations of bright multimode twin-beams using feed-forward;

• MDR is suggested to characterize the statistical properties of multimode beams after feed-forward;

• It is shown theoretically and experimentally that photon-number fluctuations of multimode twin-beams can be reduced without reduction of correlations between the beams.

Acknowledgements





Further plans and the follow-up

- Homodyning on the bright multimode squeezed states / twin beams
- Homodyning without local oscillator Laszlo Ruppert

Thank you for attention!

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