Narrowing the temporal wavepacket of SPDC photons

<u>Mikołaj Lasota</u>, Karolina Sędziak, Piotr Kolenderski

Faculty of Physics, Astronomy and Informatics, Nicolaus Copernicus University, Toruń, Poland



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NATIONAL LABORATORY OF ATOMIC, MOLECULAR AND OPTICAL PHYSICS

Background

- Major problem of realistic quantum communication: errors
- An efficient way to reduce errors independent from the real signals: temporal filtering



 Main limitation on temporal filtering in long-distance communication: temporal broadening of signal



Detection scheme for SPDC photons



- Our goal: investigating the connection between the temporal width of a photon arriving at Alice's detector and the type of spectral correlation between this photon and the corresponding photon travelling to Bob
- Temporal width of a given photon = standard deviation of the probability distribution function for the arrival time of this photon to the detector

Types of spectral correlation:
negative (typical)



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 - negative (typical)
 - no correlation



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• The analytical approximation of the spectral wave function (at the input of the fibers):

$$\phi(\nu_A, \nu_B) = \frac{1}{\sqrt{\pi}\sqrt{\sigma_A \sigma_B \sqrt{1 - \rho^2}}} \exp\left(-\frac{1}{2\left(1 - \rho^2\right)} \left(\frac{\nu_A^2}{\sigma_A^2} + \frac{\nu_B^2}{\sigma_B^2} - \frac{2\nu_A \nu_B \rho}{\sigma_A \sigma_B}\right)\right)$$

where

 σ_A , σ_B - spectral widths of the emitted photons

ho – spectral correlation coefficient (negative: -1 <
ho < 0, no correlation: ho = 0 , positive: 0 <
ho < 1)

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 - negative (typical)
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• The analytical approximation of the spectral wave function (at the input of the fibers) – for symmetric source:

$$\phi(\nu_A, \nu_B) = \frac{1}{\sigma_0 \sqrt{\pi} \sqrt[4]{1 - \rho^2}} \exp\left(-\frac{\nu_A^2 + \nu_B^2 - 2\nu_A \nu_B \rho}{2\sigma_0^2 (1 - \rho^2)}\right)$$

where

- $\sigma_0 {
 m spectral}$ width of the emitted photons
- ρ spectral correlation coefficient (negative: $-1<\rho<0,$ no correlation: $\rho=0$, positive: $0<\rho<1)$

 Initial temporal correlation: opposite to the spectral correlation



 Initial temporal correlation: opposite to the spectral correlation



• Temporal wave function at the output of the fibers of length *L*:

$$\psi_{L_A L_B}(t_A, t_B) = \int \mathrm{d}t'_A \,\mathrm{d}t'_B \,\mathcal{S}_A(t_A, t'_A, L_A) \mathcal{S}_B(t_B, t'_B, L_B) \psi(t'_A, t'_B)$$

where

$$\psi(t_A, t_B)$$
 - initial temporal wave function,
 $S_k(t_k, t'_k, L_k) = \frac{1}{\sqrt{4\pi i \beta_k L_k}} \exp\left(\frac{i(t_k - t'_k)^2}{4\beta_k L_k}\right)$
 β_k – group velocity dispersion

 Initial temporal correlation: opposite to the spectral correlation



• Effect of the propagation:



• Paerson correlation coefficient:

$$r_{t_A t_B} = \frac{E\left[(t_A - E[t_A])(t_B - E[t_B])\right]}{\sqrt{\left[E\left[(t_A - E[t_A])^2\right]E\left[(t_B - E[t_B])^2\right]}}$$

• For the wavefunction $\psi_L(t_A, t_B)$:

$$r_{t_A t_B}(L,\rho) = -\rho \frac{1 - 4\sigma_0^4 \beta^2 L^2 (1-\rho^2)}{1 + 4\sigma_0^4 \beta^2 L^2 (1-\rho^2)}$$



• The Pearson coefficient as a function of the propagation distance and the spectral correlation coefficient. The values $\beta = -1.15 \times 10^{-26} \frac{s^2}{m}$ and $\sigma_0 = 1.57$ THz are assumed.

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• The distance at which $r_{t_A t_B} = 0$:

$$L_0 = \frac{1}{2\sigma_0^2\beta\sqrt{1-\rho^2}}$$

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- Temporal width of a given photon = standard deviation of the probability distribution function for the arrival time of this photon to the detector

• Temporal width of Alice's photon if the **global time reference** is **available** but the **detection time of Bob's photon** is **not known**:

$$au_A(\sigma_A) = \sqrt{rac{g(x_A^2)}{2\sigma_A^2(1-
ho^2)}}$$
 where $g(x) = 1 + x(1-
ho^2)$ and $x_Y = 2\sigma_Y^2 \beta_Y L_Y$



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 Temporal width of Alice's photon if the global time reference is available and the detection time of Bob's photon is known:

$$\tau_{Ah}(\sigma_A, \sigma_B) = \sqrt{\frac{\left[g(-x_A x_B)\right]^2 + (x_A + x_B)^2}{2\sigma_A^2 g(x_B^2)}}$$



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• Temporal width of Alice's photon if the **global time reference** is **not available**, but **she can measure** $\Delta t = t_A - t_B$:

$$\tau_{Ah,\Delta t}(\sigma_A,\sigma_B) = \sqrt{\frac{\{g(x_A^2)\sigma_B^2 + g(x_B^2)\sigma_A^2 + 2g(-x_A x_B)\sigma_A \sigma_B \rho\} \left\{ [g(-x_A x_B)]^2 + (x_A + x_B)^2 \right\}}{2\sigma_A^2 \sigma_B^2 \left[g(x_A^2)g(x_B^2) - [g(-x_A x_B)]^2 \rho^2 \right]}}$$



Discrete-variable QKD setup with source of photons located in the middle:
 Alice
 PBS
 SMF fiber
 SMF fiber



• Discrete-variable QKD setup with source of photons located in the middle:



- Assumptions:
 - BB84 protocol
 - duration time of a single detection window = six temporal widths
 (99,73% probability for successful detection)

• Discrete-variable QKD setup with source of photons located in the middle:



- Technical assumptions:
 - the only source of errors: dark counts (rate d = 1kHz)
 - transmittance of the fibers: $\,T=10^{-lpha L/10}$, where $\,lpha=0.2\,{
 m dB/km}$
 - group velocity dispersion (GVD) for the SMF fibers: $2\beta_A = 2\beta_B = -2.3 \times 10^{-26} \text{s}^2/\text{m}$
 - no other setup imperfections



- Spectral widths: $\sigma_A = \sigma_B =$ **1.57 THz**
- Spectral correlation: $\rho = -0.9$ (red), $\rho = 0$ (green), $\rho = 0.9$ (blue)
- Availability of global time reference: yes (dashed), no (solid)

- Conclusions:
 - strong spectral correlation (positive or negative) is better than no correlation at all for the case with global time reference
 - positive spectral correlation is the best and negative spectral correlation is the worst for the case without global time reference

- Conclusions:
 - strong spectral correlation (positive or negative) is better than no correlation at all for the case with global time reference
 - positive spectral correlation is the best and negative spectral correlation is the worst for the case without global time reference
 - for spectrally narrow pairs the situation can be totally different!

Symmetric quantum key distribution



- Spectral widths: $\sigma_A = \sigma_B = 10 \text{ GHz}$
- Spectral correlation: $\rho = -0.9$ (red), $\rho = 0$ (green), $\rho = 0.9$ (blue)
- Availability of global time reference: yes (dashed), no (solid)

SPDC spectral wavefunction in another parametrization

• Spectral wavefunction at the input of the fibers – "old parametrization":

$$\phi(\nu_A, \nu_B) = \frac{1}{\sigma_0 \sqrt{\pi} \sqrt[4]{1 - \rho^2}} \exp\left(-\frac{\nu_A^2 + \nu_B^2 - 2\nu_A \nu_B \rho}{2\sigma_0^2 (1 - \rho^2)}\right)$$

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 Spectral wavefunction at the input of the fibers – "new parametrization" (collinear case):

$$\phi(\nu_A,\nu_B) = M \exp\left(-\frac{(\nu_A - \nu_B)^2}{\sigma^2} - \frac{(\nu_A + \nu_B)^2 \tau_p^2}{4}\right)$$

where

 σ – characteristic width of the effective phase-matching function τ_p – pump laser pulse duration

- Temporal width of Alice's photon in the new parametrization:
 - when Alice knows only the emission time of a given pair of photons

$$\tau_A(\sigma,\tau_p) = \frac{\sqrt{\left(\tau_p^2 + \beta_A^2 L_A^2 \sigma^2\right) \left(4 + \sigma^2 \tau_p^2\right)}}{2\sigma \tau_p}$$

- when Alice knows both the emission time of a given pair of photons and the detection time of Bob's photon

$$\tau_{Ah}(\sigma,\tau_{p}) = \sqrt{\frac{16\left(\tau_{p}^{2} - \beta_{A}L_{A}\beta_{B}L_{B}\sigma^{2}\right)^{2} + \left(\beta_{A}L_{A} + \beta_{B}L_{B}\right)^{2}\left(4 + \sigma^{2}\tau_{p}^{2}\right)^{2}}{4\left(4 + \sigma^{2}\tau_{p}^{2}\right)\left(\tau_{p}^{2} - \beta_{A}L_{A}\beta_{B}L_{B}\sigma^{2}\right)}}$$

- when Alice knows only the detection time of Bob's photon

$$\tau_{Ah,\Delta t}(\boldsymbol{\sigma},\tau_p) = \frac{\sqrt{\left[16 + \left(\beta_A L_A + \beta_B L_B\right)^2 \boldsymbol{\sigma}^4\right] \tau_p^2 + 4\boldsymbol{\sigma}^2 \left(\beta_A L_A - \beta_B L_B\right)^2}}{2\boldsymbol{\sigma}\tau_p}$$



• Minimization of Alice's photon temporal width in the case of symmetric setup (two SMFs of 10 km length, $\sigma = 3.25$ THz):

• Minimization of Alice's photon temporal width in the case of symmetric setup ($\beta_A L_A = \beta_B L_B \equiv \beta L$):

-
$$\tau_A^{\min}(\sigma) = \frac{|\beta|L\sigma^2 + 2}{2\sigma}$$
 for $\tau_p^{\text{opt}} = \sqrt{2|\beta|L}$
- $\tau_{Ah}^{\min}(\sigma) = \frac{2\sqrt{|\beta|L(\beta^2L^2\sigma^4 + 4)}}{|\beta|L\sigma^2 + 2}$ for $\tau_p^{\text{opt}} = \sqrt{2|\beta|L}$

$$- \tau_{Ah,\Delta t}(\sigma) = \frac{\sqrt{\beta^2 L^2 \sigma^4 + 4}}{\sigma}$$

independent from au_p

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$$\tau_{Ah,\Delta t}(\sigma) = \frac{\sqrt{\beta^{2}L^{2}\sigma^{4}+4}}{\sigma} \quad \text{independent from } \tau_{p}$$

• Optimal value of spectral correlation coefficient for the symmetric setup:

$$\rho^{opt} = \frac{2 - |\beta| L \sigma^2}{2 + |\beta| L \sigma^2}$$

• Absolute minimum of Alice's photon temporal width in the case of symmetric setup:

-
$$\tau_A^{abs} = \tau_{Ah}^{abs} = \sqrt{2|\beta|L}$$
 for $\tau_p^{opt} = \sqrt{2|\beta|L}$ and $\sigma^{opt} = \sqrt{2/|\beta|L}$

-
$$au_{Ah,\Delta t}^{
m abs} = 2\sqrt{|eta|L}$$
 for $\sigma^{
m opt} = \sqrt{2/|eta|L}$ and arbitrary au_p

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m opt}=\sqrt{2/|eta|L}$ and arbitrary au_p

 Optimal type of spectral correlation between SPDC photons: no correlation



• $\log_{10} \tau_A$ plotted for symmetric setup with two 1km-long SMFs



• $\log_{10} \tau_A$ plotted for symmetric setup with two 1km-long SMFs



• $\log_{10} \tau_{Ah}$ plotted for symmetric setup with two 1km-long SMFs



• $\log_{10} \tau_{Ah,\Delta t}$ plotted for symmetric setup with two 1km-long SMFs

• Discrete-variable QKD setup with source of photons located in the middle:



- Assumptions:
 - BB84 protocol
 - optimized duration time of a single detection window

Influence of the duration time of detection windows



• Key generation rate for the asymmetric QKD setup plotted for the following duration time of the detection window: $12\tau_{Bh}$ (red), $6\tau_{Bh}$ (orange), $3\tau_{Bh}$ (green), $1.5\tau_{Bh}$ (blue), $0.6\tau_{Bh}$ (black)

Potential to improve QKD security



• Logarithm of secure key rate calculated for symmetric QKD setup with two SMF fibers and for optimized duration time of a single detection window $\sigma = 3.25$ THz, $\tau_p = 0.1$ ps - $\sigma = 3.25$ THz, $\tau_p = \sqrt{2|\beta|L}$ - $\sigma = \sqrt{2/|\beta|L}$, $\tau_p = \sqrt{2|\beta|L}$

Our papers

- 1. K. Sedziak, M. Lasota, P. Kolenderski, Optica 4, 84-89 (2017):
 - investigation of the dependence of temporal width of SPDC photons propagated in SMF fibers on the type of spectral correlation between them
 - example of application: symmetric QKD setup
- 2. M. Lasota, P. Kolenderski, arXiv:1702.05165 (2017):
 - consideration of asymmetric QKD setup
 - investigation of the possibility of extending QKD security by adding chromatic dispersion to the shorter fiber
- 3. K. Sedziak, M. Lasota, P. Kolenderski, arXiv:1711.06131 (2017):
 - optimization of the pump laser for a given SPDC source for QC applications
 - designing optimal SPDC source for QC applications
 - experimental results regarding wavepacket narrowing of SPDC photons propagated in SMF fibers



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