

Quantum optics with photons, phonons and plasmons



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Shailesh Kumar, Christian Kothe, Mikael Lassen, Niels Israelsen,
Adriano Berni, Mario Usuga and Ulrik L. Andersen

Department of Physics, Technical University of Denmark

Some recent collaborators:

Uni Palacky (R. Filip, L. Mista, V. Usenko M. Jezek, J. Fiurasek, P. Marek)

Uni Erlangen (G. Leuchs, Ch. Marquardt, Ch. Wittmann, Ch. Muller, Ch. Gabriel...)

Uni Queensland (W. Bowen, G. Harris, T. Ralph)

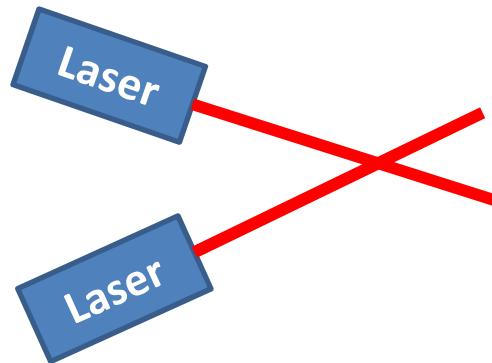
Kista (G. Bjork)

Uni. Copenhagen (A. Sørensen, P. Lodahl, E. Polzik)

Brussel Uni (N. Cerf, J. Niset)



It is good to have LARGE non-linearities



Measurement induced nonlinearities -> Linear optical computing

Not resource efficient!

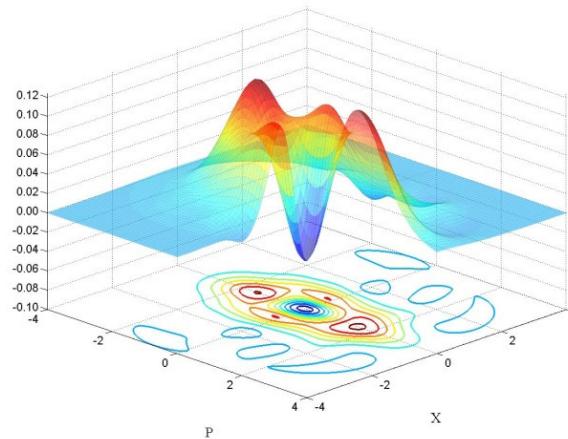
Giant non-linearity between photons -> Deterministic quantum logic gates

Large nonlinearity between photons -> Near-Deterministic quantum logic gates

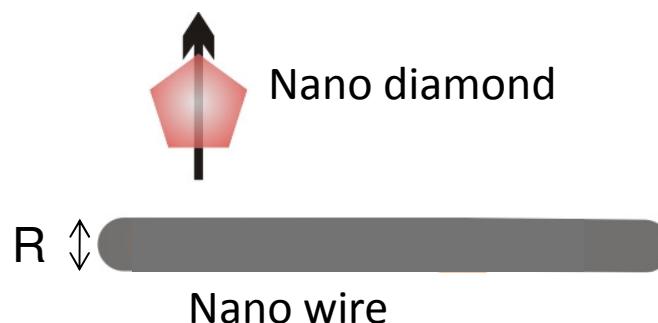
Giant non-linearity -> Non-Gaussian states

Outline

Measurement induced nonlinearity



Potential candidate for deterministic large nonlinearities.



Measurement induced giant nonlinearity: Implementation of a Hadamard gate for coherent states

A. Tipsmark¹, R. Dong¹, M. Jezek^{1,2}, A. Laghaout¹, P. Marek², ULA¹

- 1) Technical University of Denmark (DTU)
- 2) Palacky University, Olomouc, Czech Republic

Cat state quantum computing

Ralph et al PRA 68, 042319 (2003)

Computational basis: $|\alpha\rangle, |-\alpha\rangle$

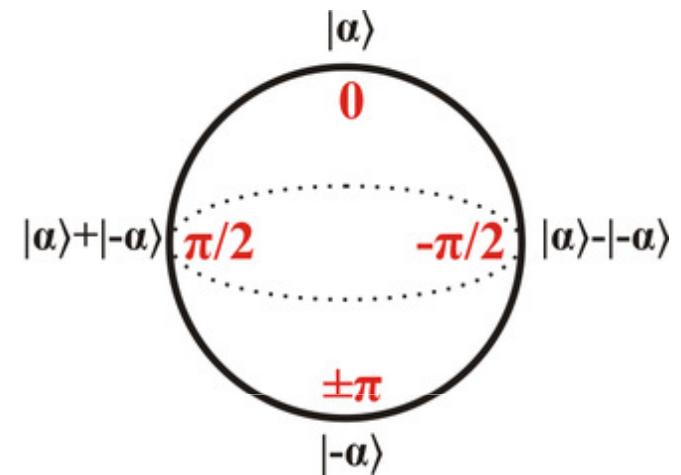
Diagonal basis: $|\alpha\rangle \pm |-\alpha\rangle$

Universal set of gates:

X-gate $|\alpha\rangle \rightarrow |-\alpha\rangle$

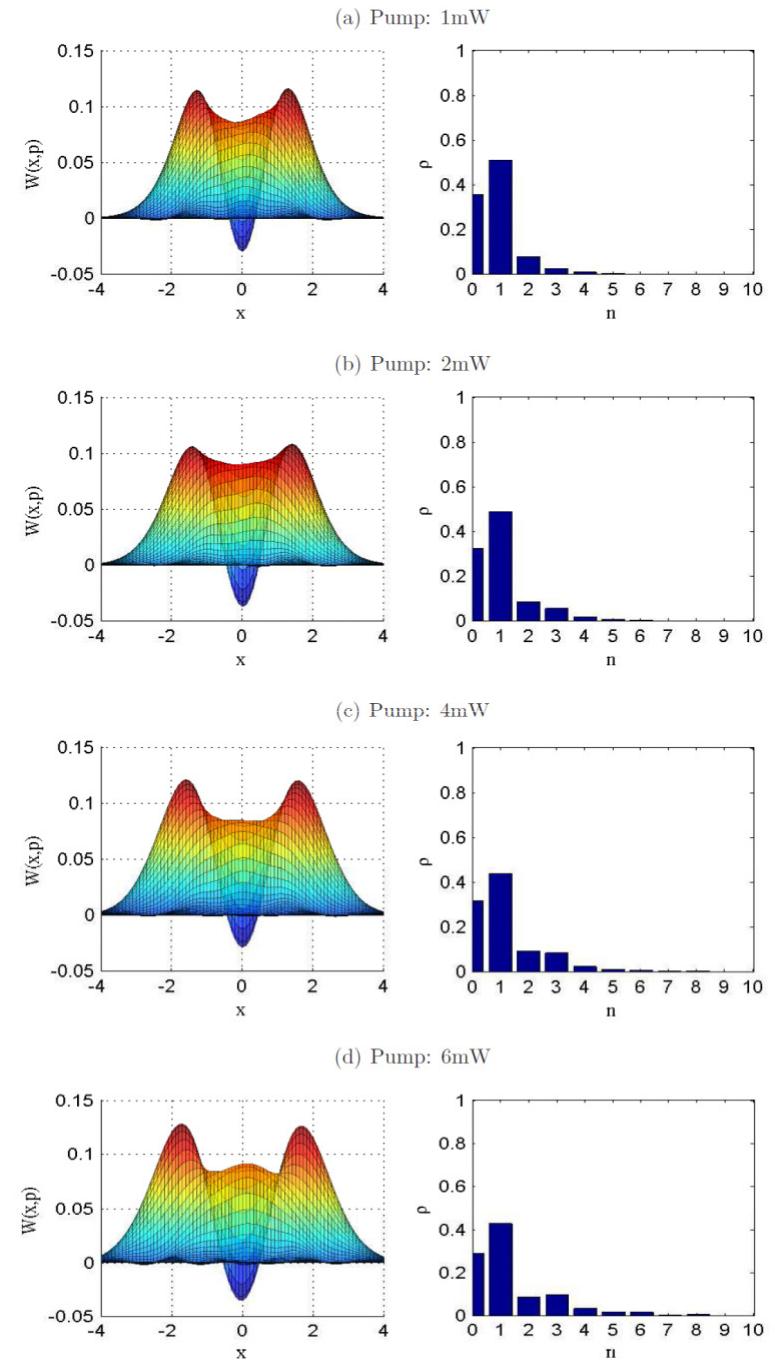
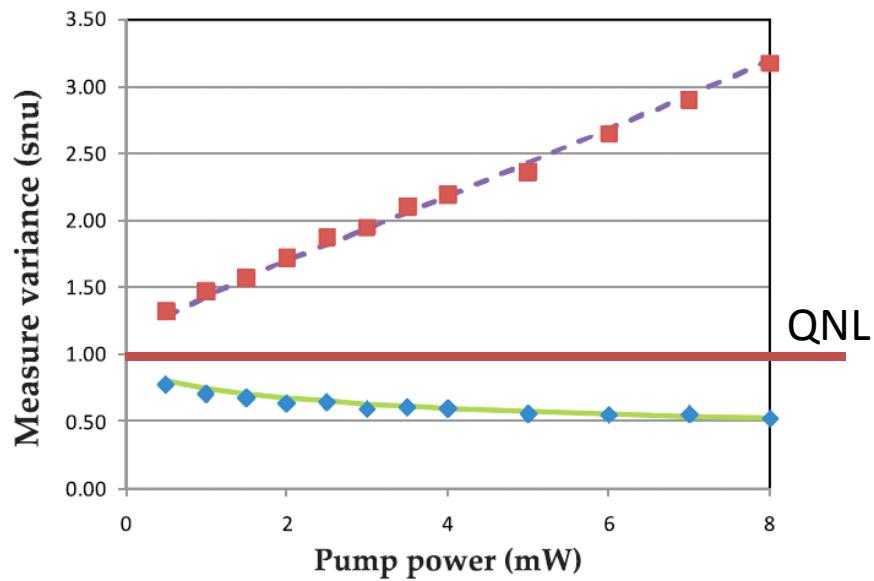
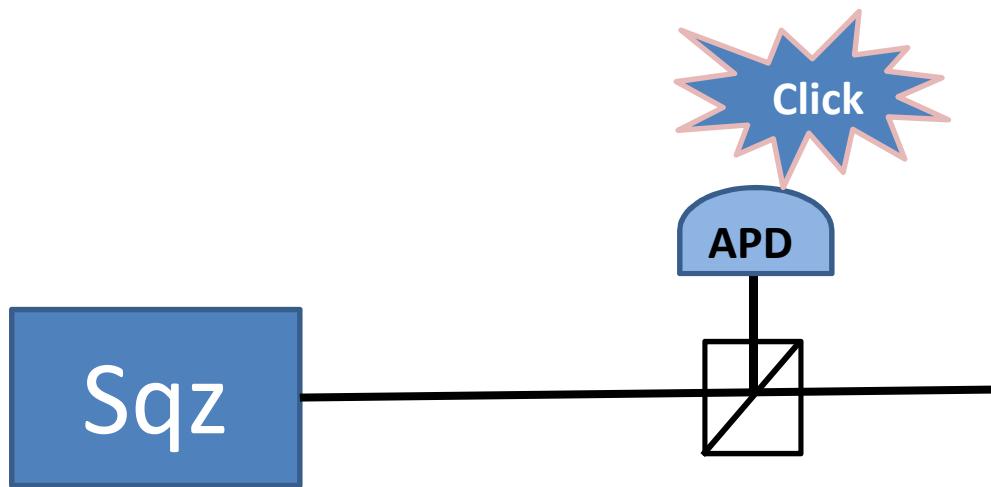
Z-rotation $|\alpha\rangle + |-\alpha\rangle \rightarrow |\alpha\rangle + e^{i\theta} |-\alpha\rangle$

Hadamard gate $|\pm\alpha\rangle \rightarrow |\alpha\rangle \pm |-\alpha\rangle$



Controlled Z-rotation

Preparation of odd cats



Ourjoumtsev et al Science 312, 83 (2006); Nielsen et al PRL 97, 083604 (2006); Wakui et al, Opt. Express 15, 3568 (2007).

Implementing the gates by teleportation

Lund, Ralph, Haselgrove PRL 100, 030503 (2008)

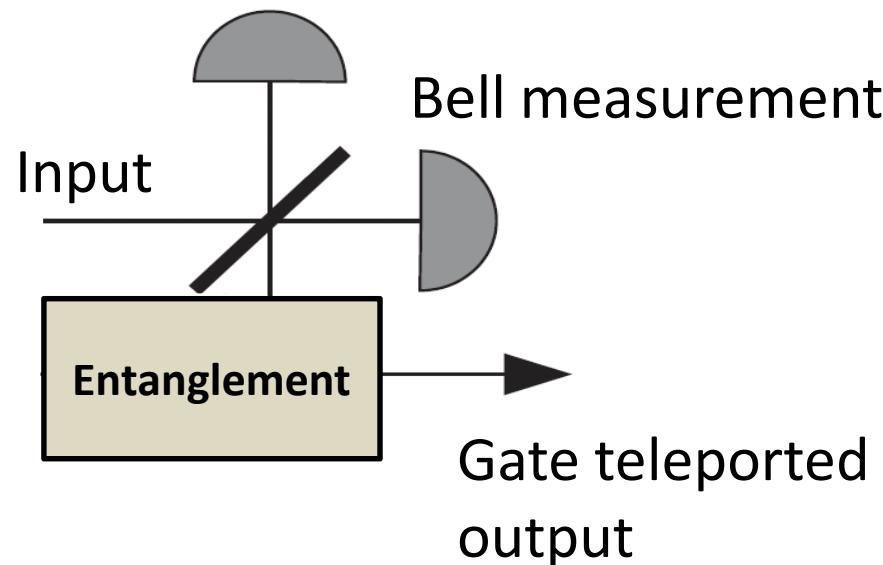
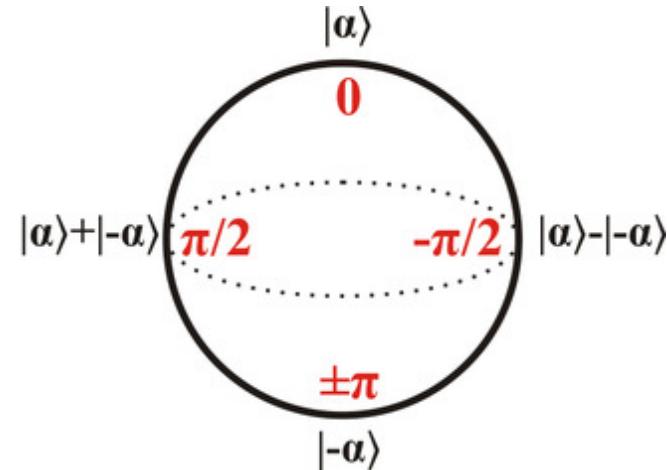
Universal set of gates:

X-gate

Z-rotation

Hadamard gate

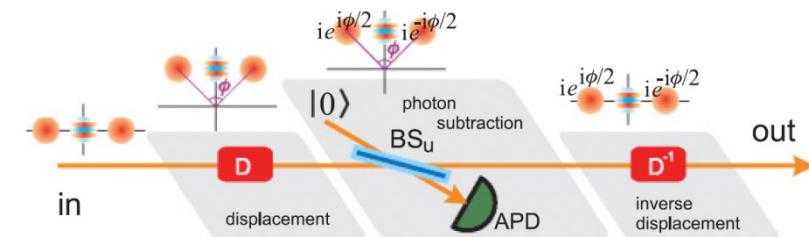
Controlled Z-rotation



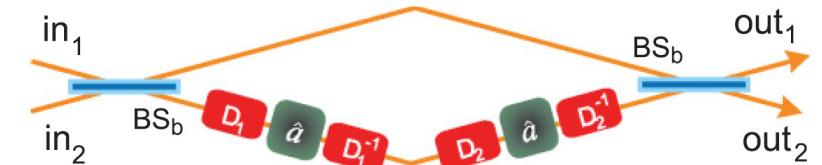
Probabilistic implementation

Marek and Fiurasek, PRA **82**, 014304 (2010)

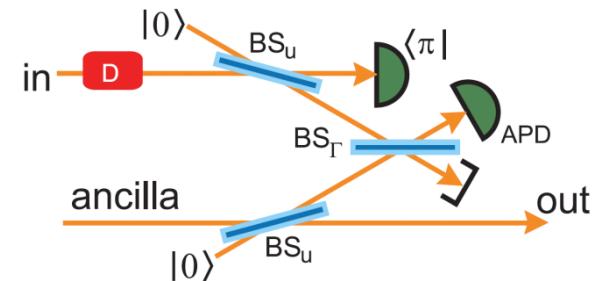
Z-rotation gate



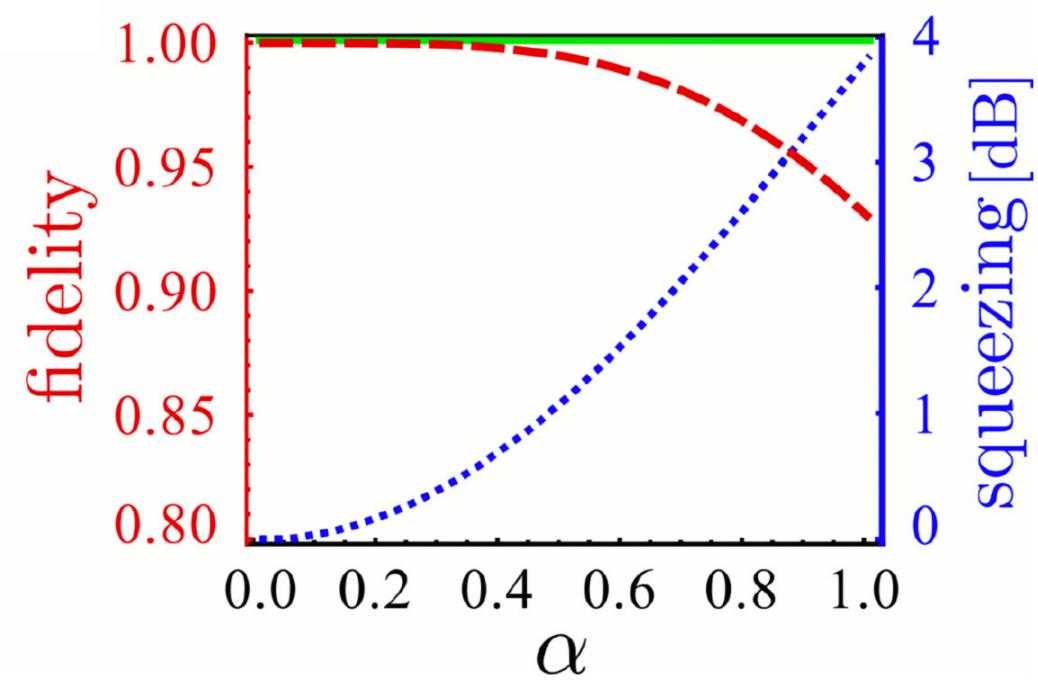
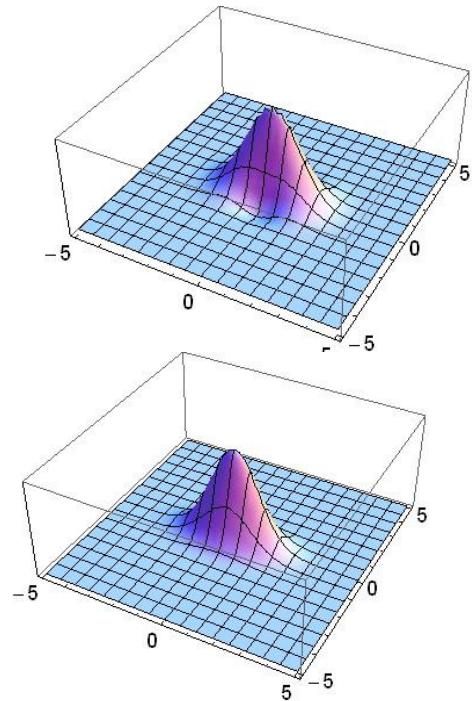
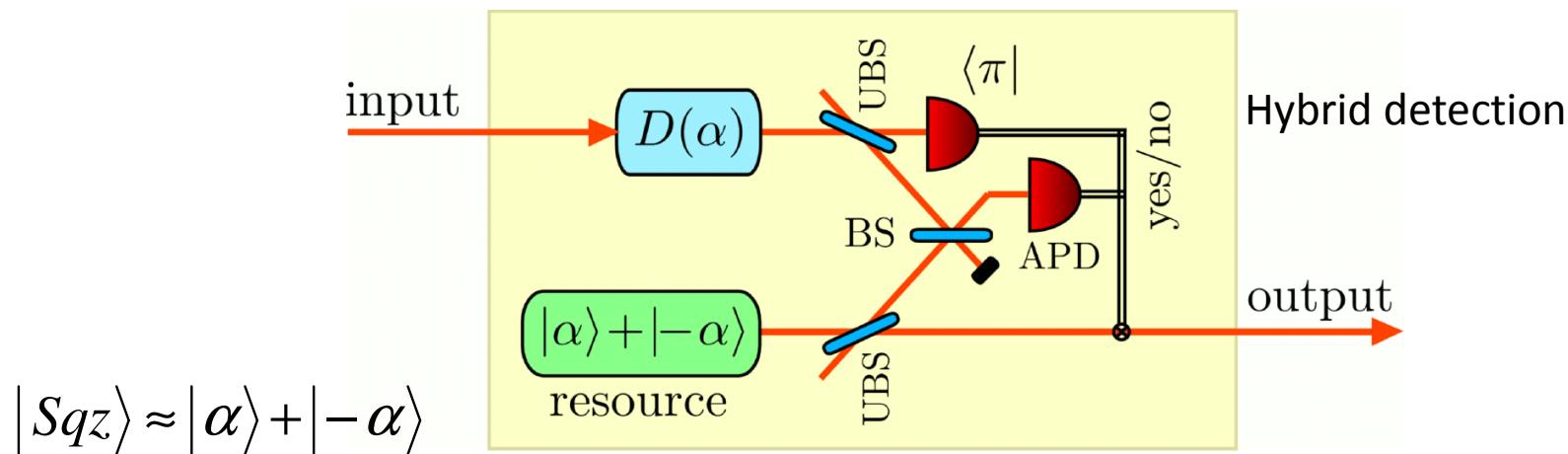
Controlled Z-rotation gate



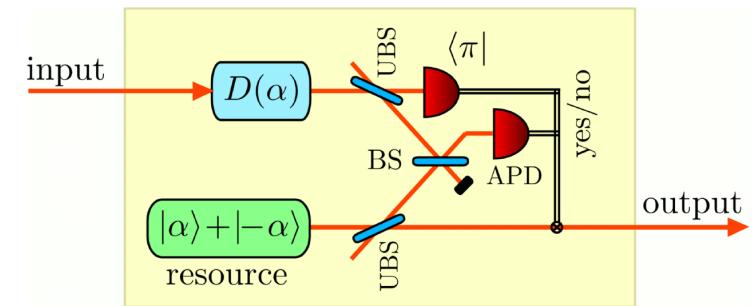
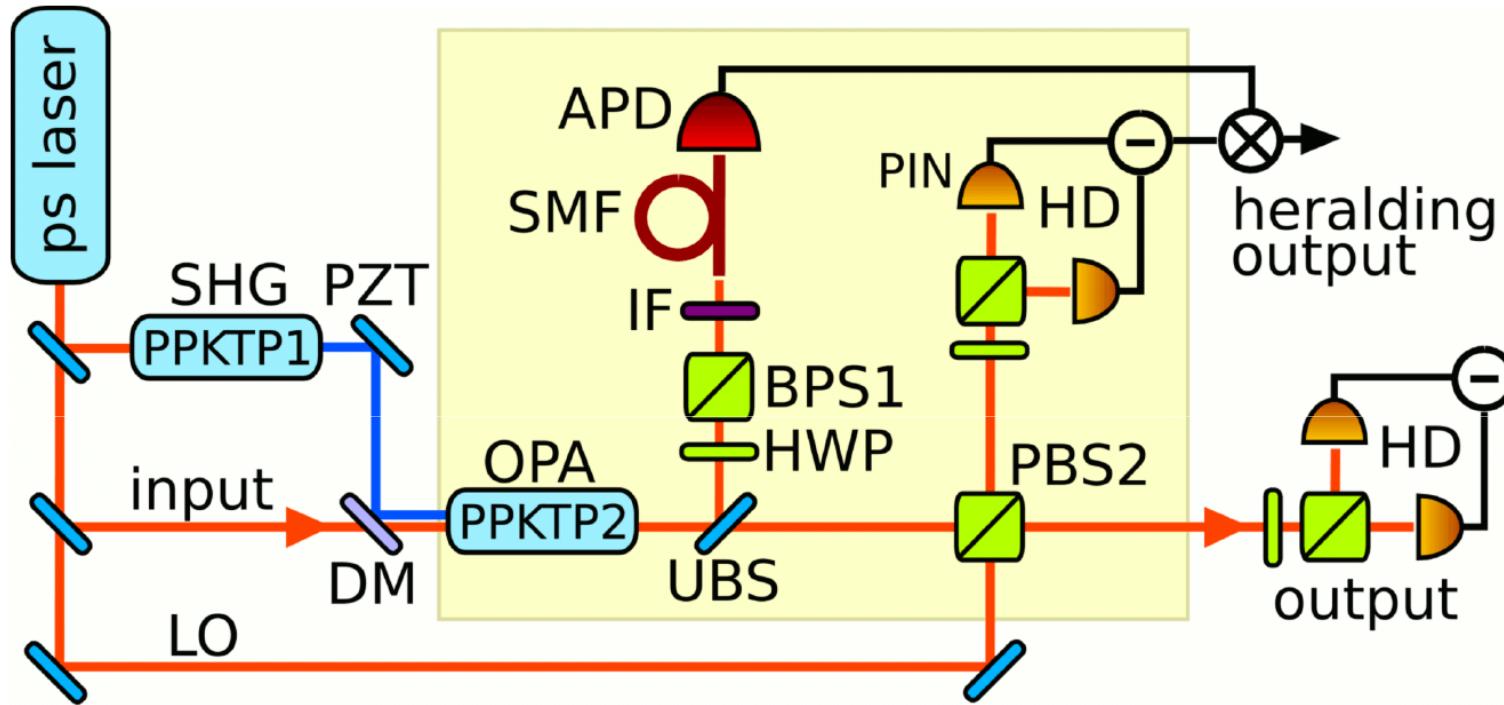
Hadamard gate



Hadamard gate $|\pm\alpha\rangle \rightarrow |\alpha\rangle \pm |-\alpha\rangle$

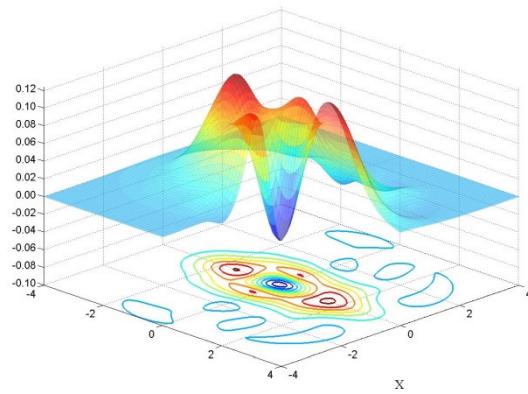


Experimental Setup



Experimental results for $\alpha = 0.8$

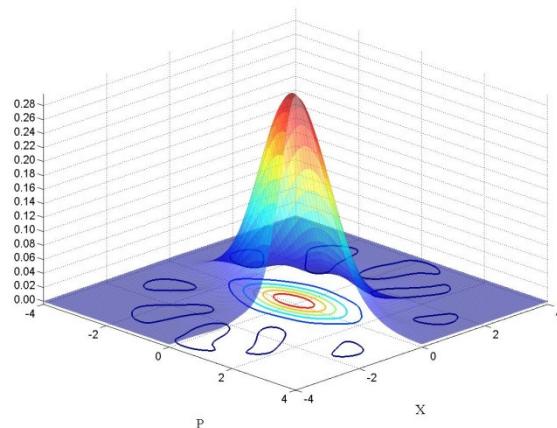
$|-\alpha\rangle \rightarrow$



$|\alpha\rangle - |-\alpha\rangle$

$F = 0.65$ for $\alpha = 0.8$

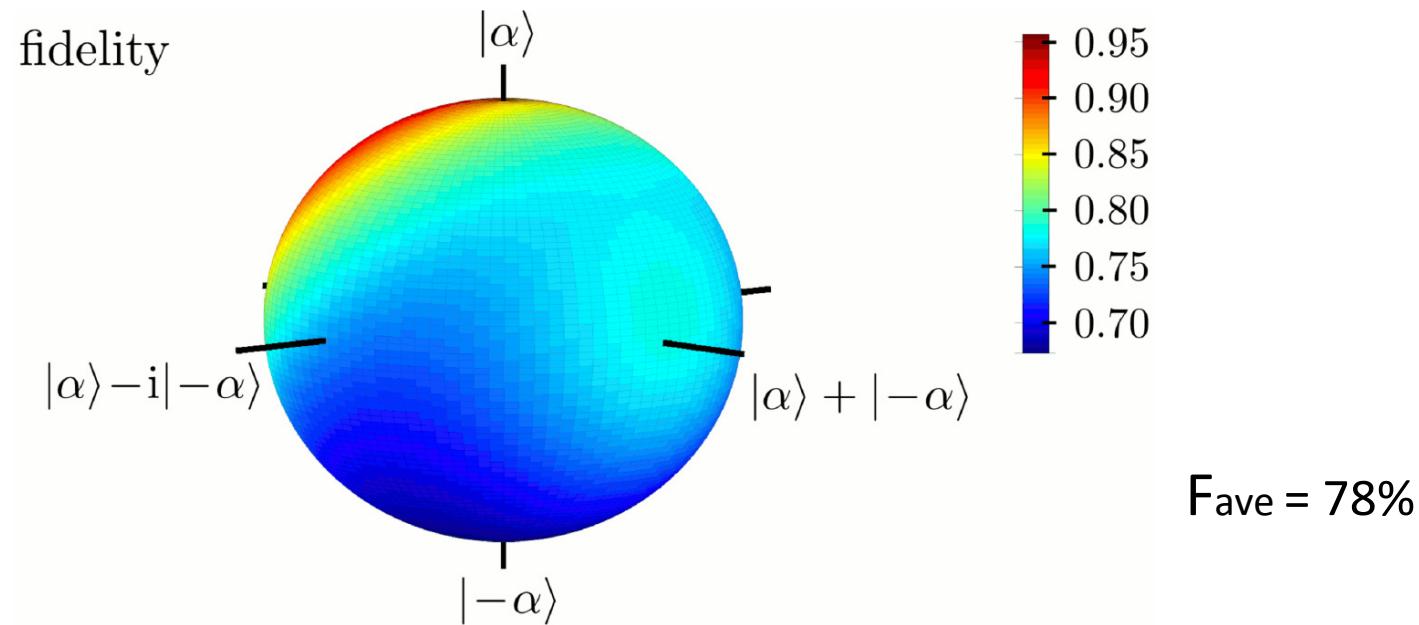
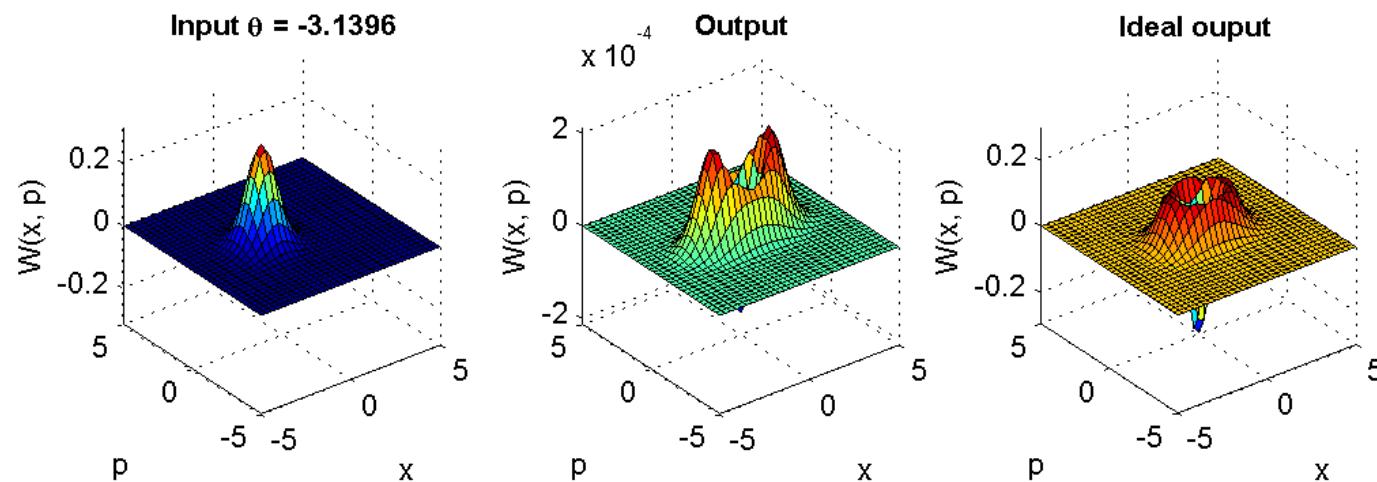
$|\alpha\rangle \rightarrow$



$|\alpha\rangle + |-\alpha\rangle$

$F = 0.94$ for $\alpha = 0.8$

Simulations based on experimental parameters



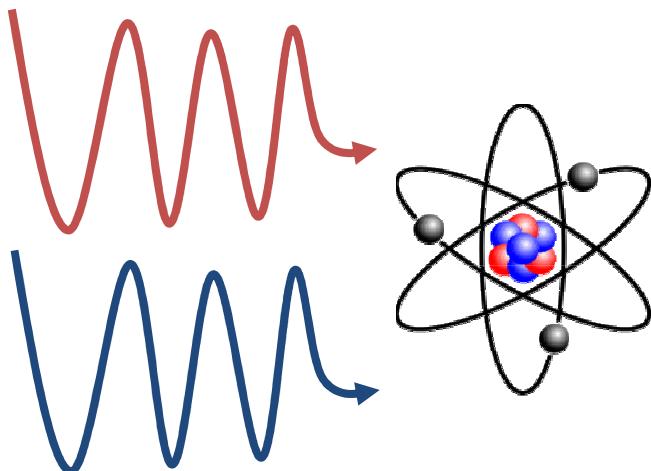
Towards non-linear interactions mediated by plasmons

Alexander Huck, Shailesh Kumar, Ying Wei Lu and Ulrik L. Andersen

Department of Physics
Technical University of Denmark



Strong light-matter coupling is what we need!



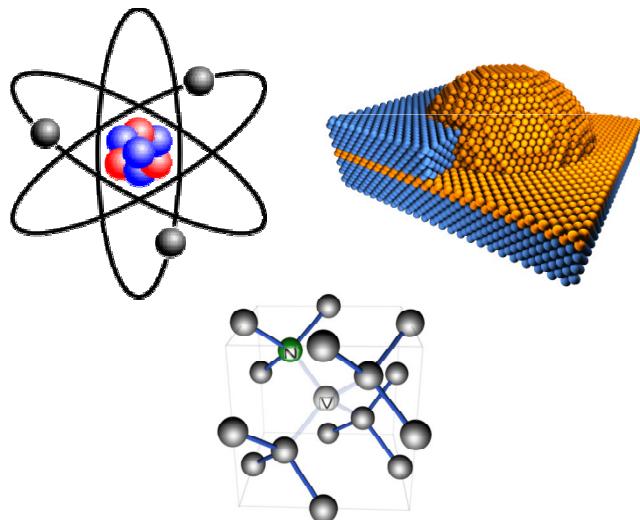
Saturation
Photon blockade
Kerr effect

Quantum computing
Quantum repeater

How is efficient coupling obtained?

Decay rate: $\Gamma \propto \mu^2 E^2 g(\omega)$

Dipole moment



Large oscillation strength!

Electric field density

$$E^2 \propto \frac{1}{V}$$

Density of states

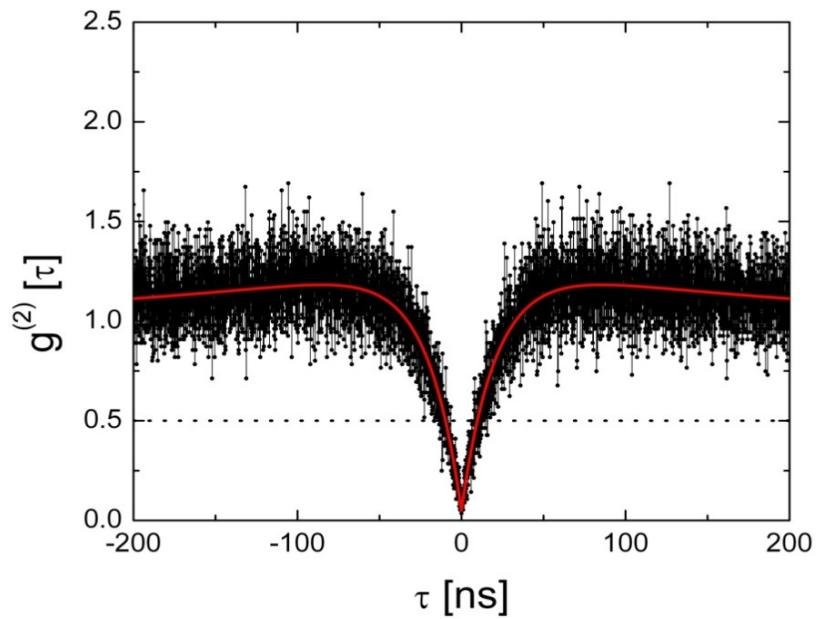
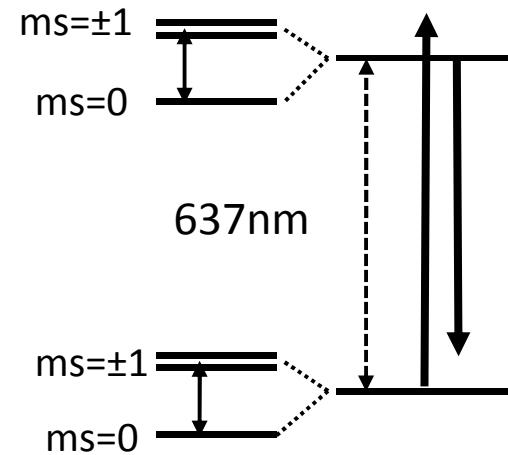
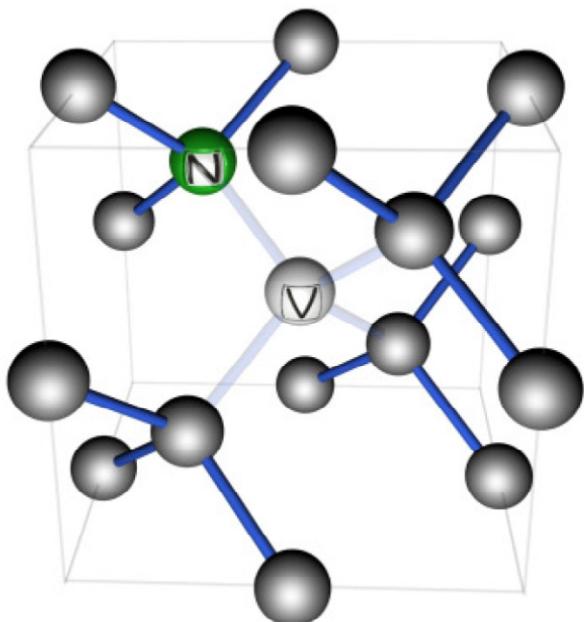
$$g(\omega) \propto \frac{1}{d\omega/dk}$$

Strong field confinement!

Small group-velocity!

Nitrogen-Vacancy center in diamond

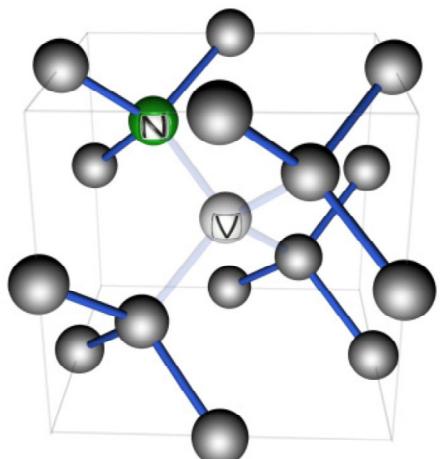
Diamond crystal with NV



How is efficient coupling obtained?

Decay rate: $\Gamma \propto \mu^2 E^2 g(\omega)$

Dipole moment



Large oscillation
strength!

Electric field
density

$$E^2 \propto \frac{1}{V}$$

Strong field
confinement!

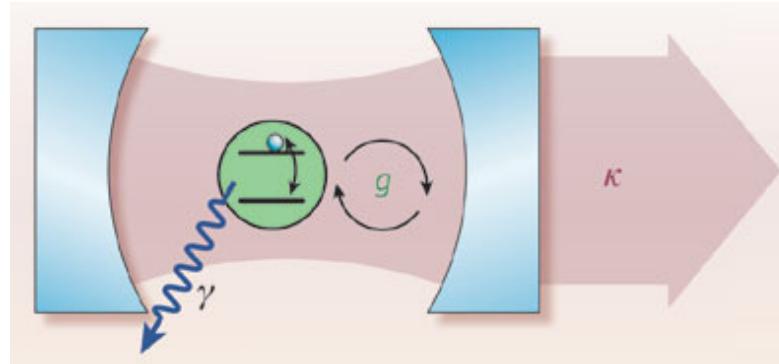
Density of states

$$g(\omega) \propto \frac{1}{d\omega/d\beta} \propto Q$$

Small group-
velocity!

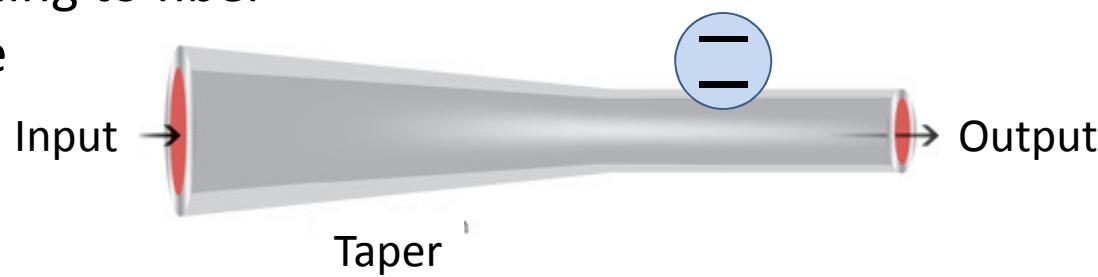
How is efficient coupling obtained?

Cavity QED



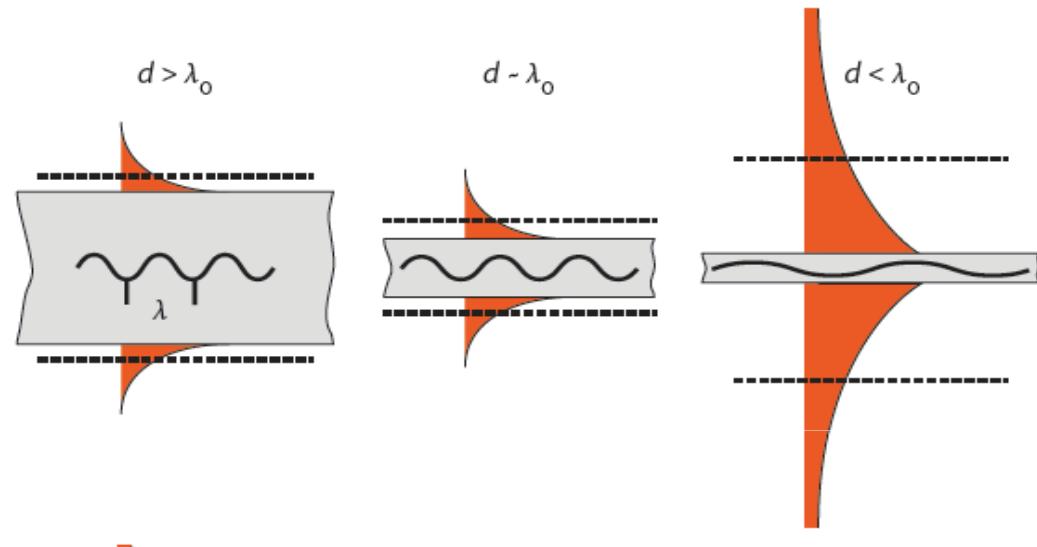
Thompson et al, PRL 68, 1132 (1992)

Coupling to fiber mode

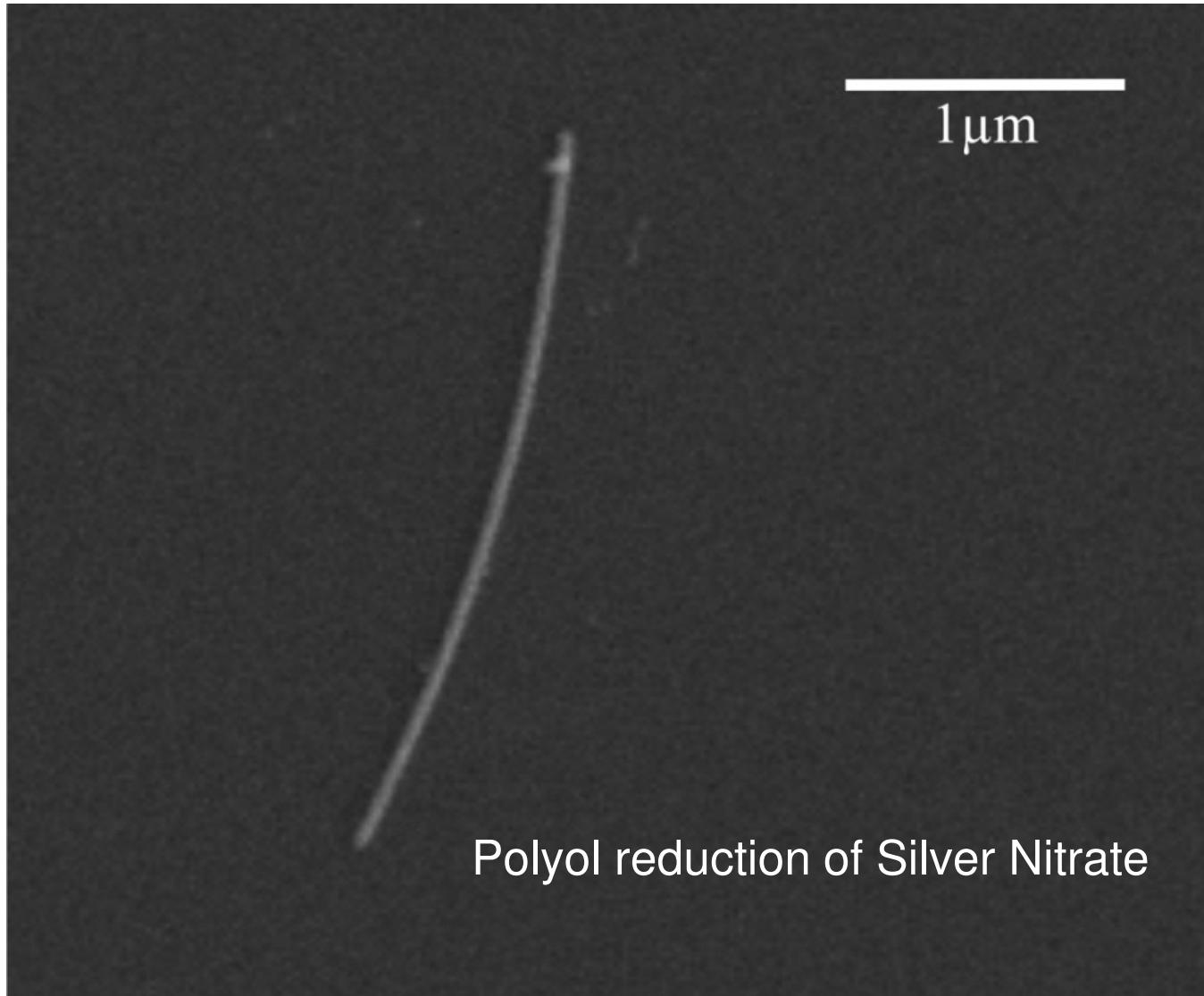


Going beyond the diffraction limit

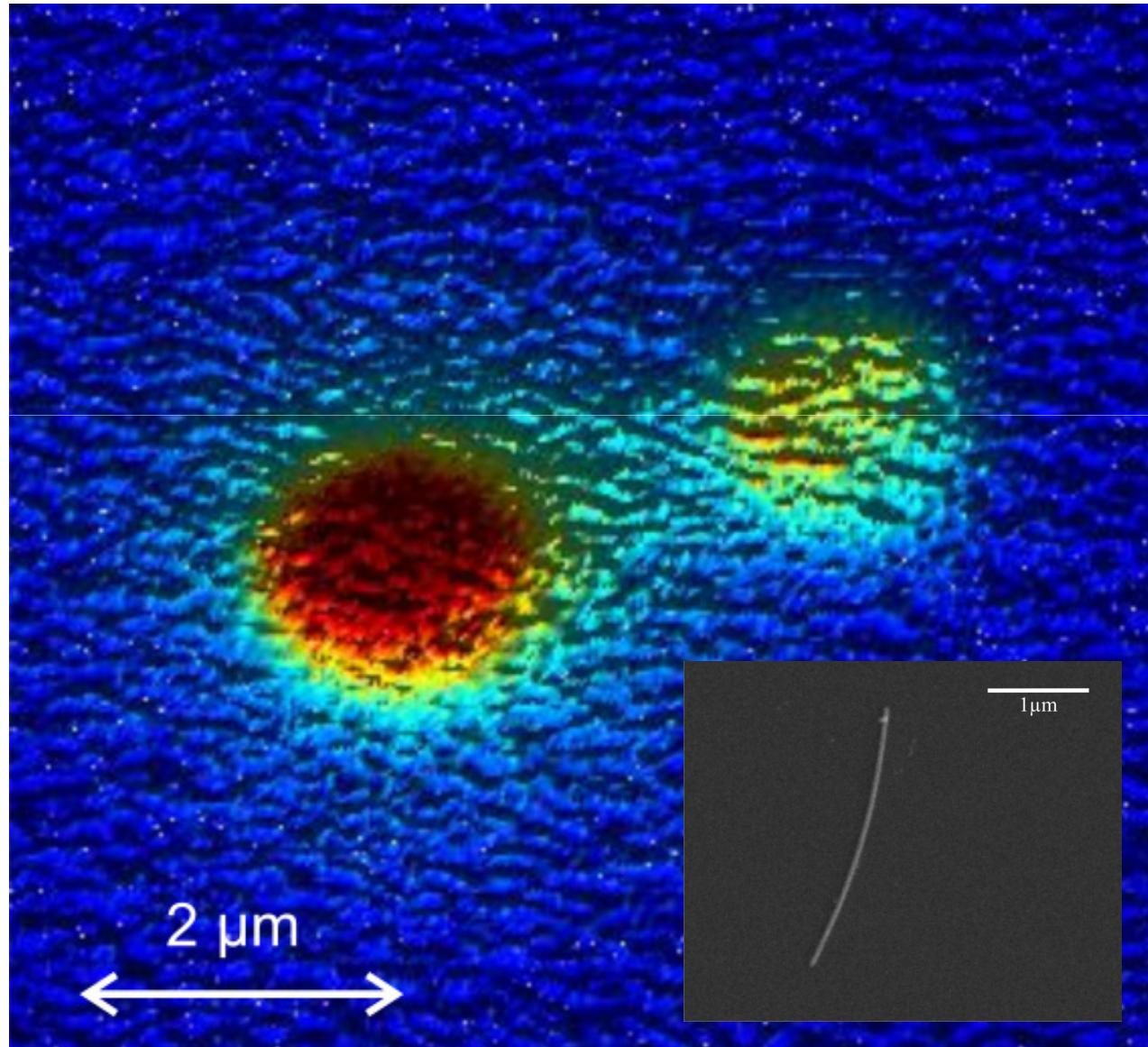
- Fiber modes:
Limited by diffraction.



Mono-crystalline silver wires

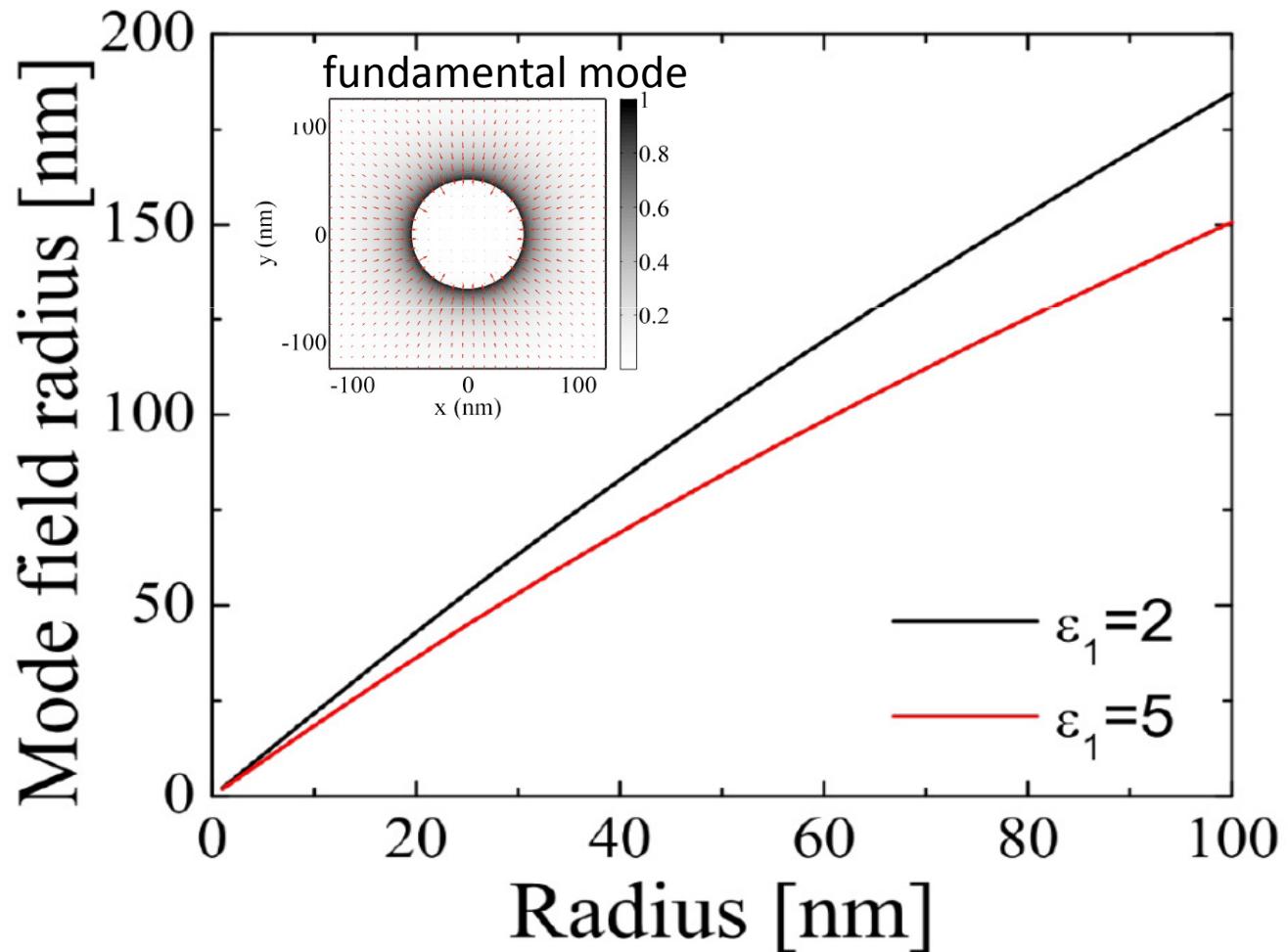


Plasmon propagation



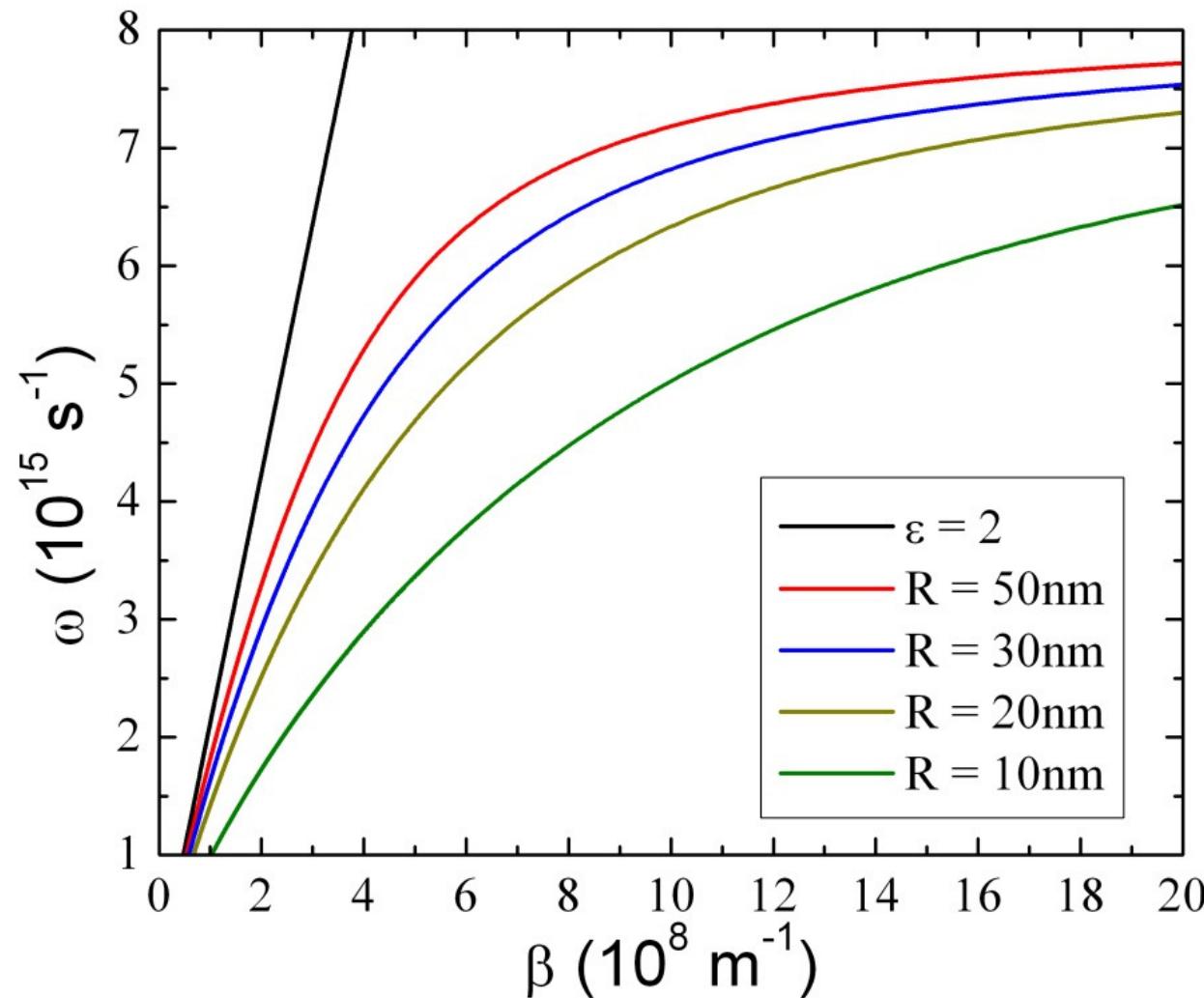
Decay rate: $\Gamma \propto \mu^2 E^2 g(\omega)$

$$E^2 \propto \frac{1}{A} \propto \frac{1}{R^2} \propto \frac{1}{(\text{Wire radius})^2}$$

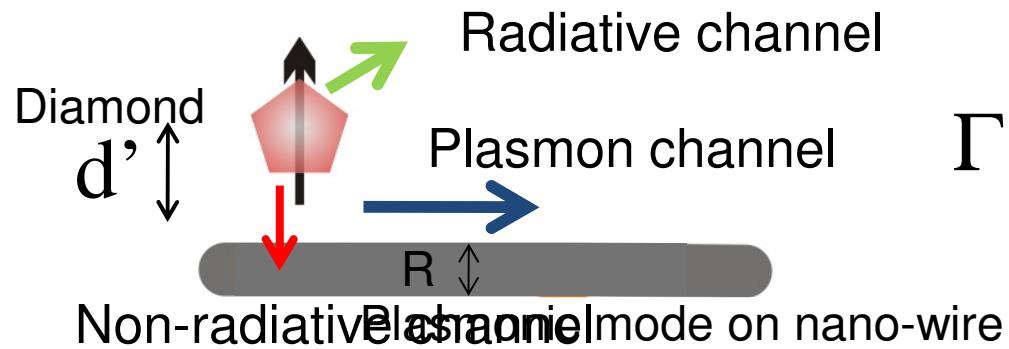


Decay rate: $\Gamma \propto \mu^2 E^2 g(\omega)$

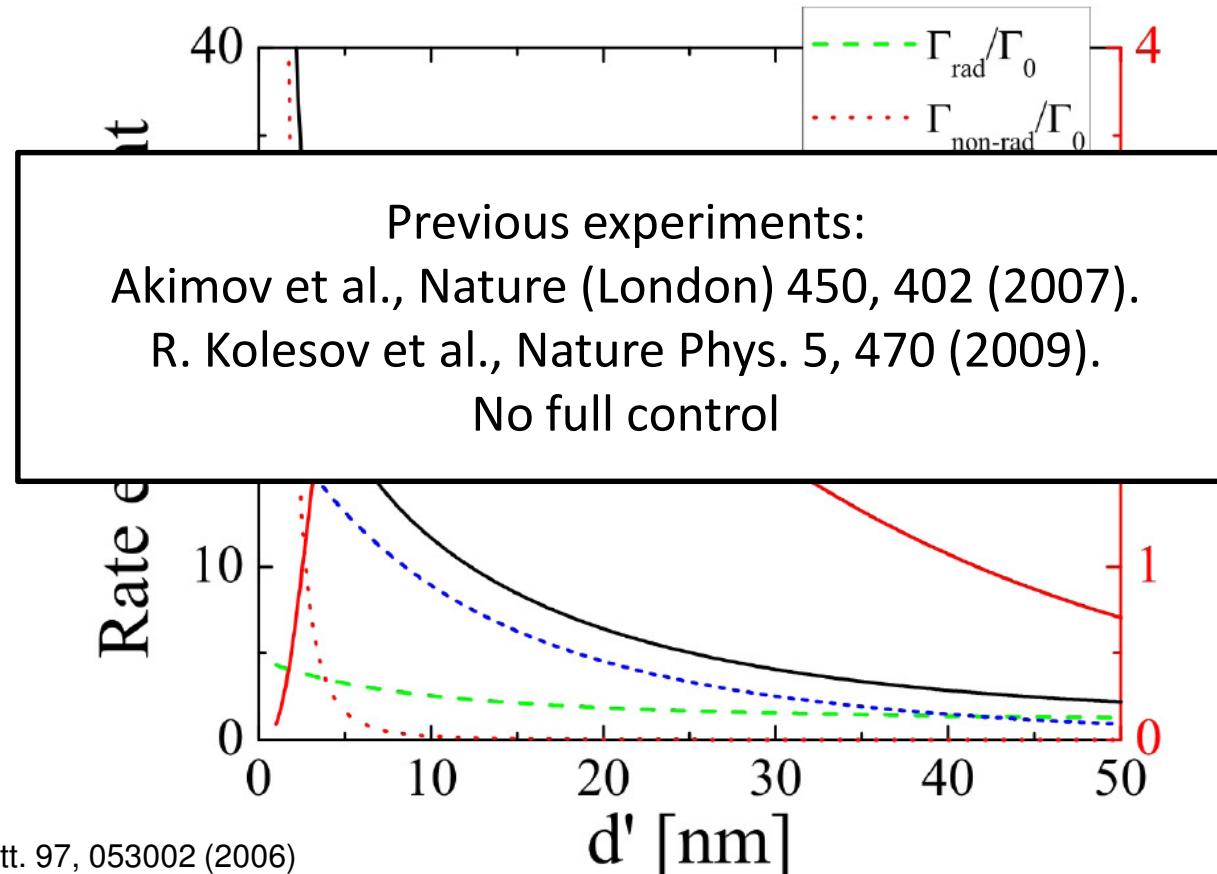
$$g(\omega) \propto \frac{1}{d\omega/d\beta} \propto \frac{1}{\text{Wire radius}}$$



Increasing the decay rate into one plasmonic mode



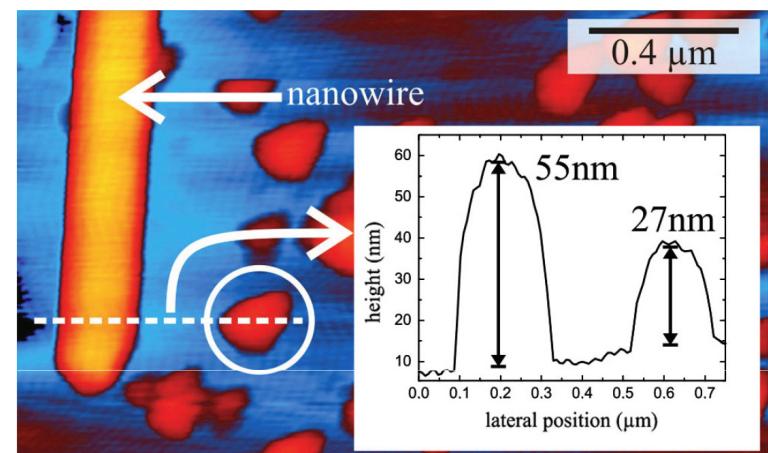
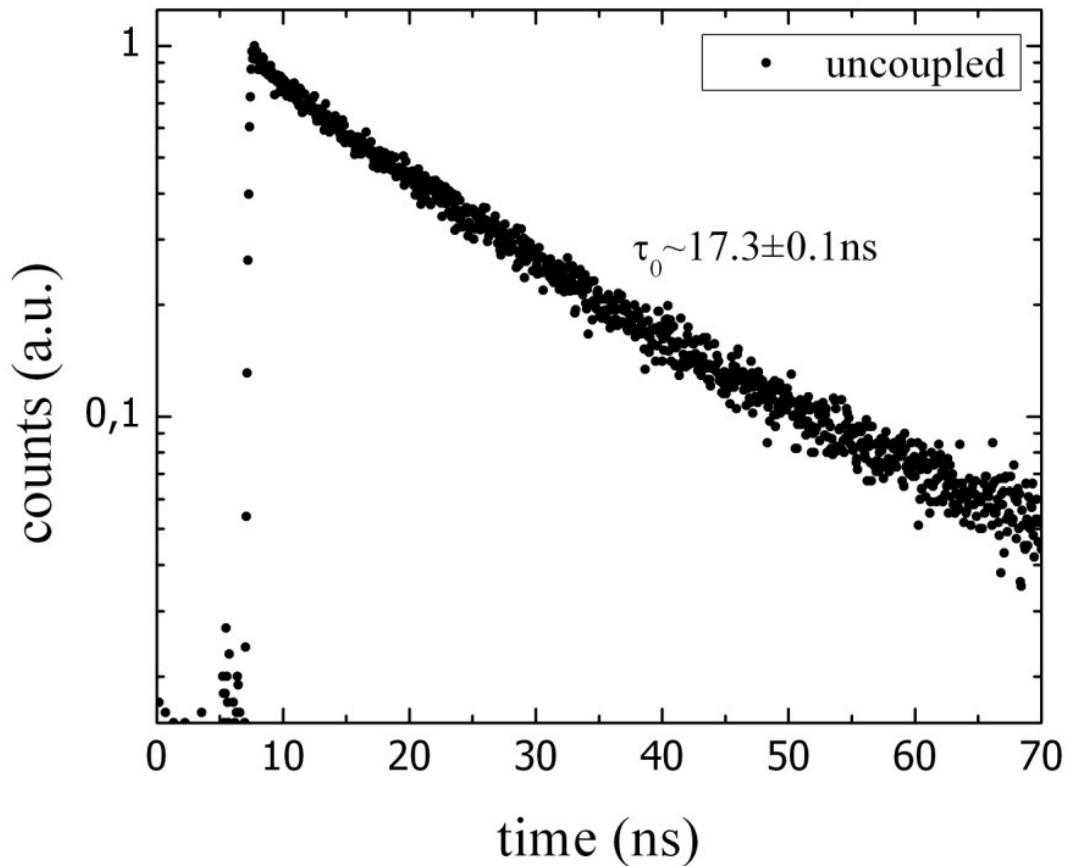
$$\Gamma^{plas} \propto \frac{\lambda}{(\text{Wire radius})^3}$$



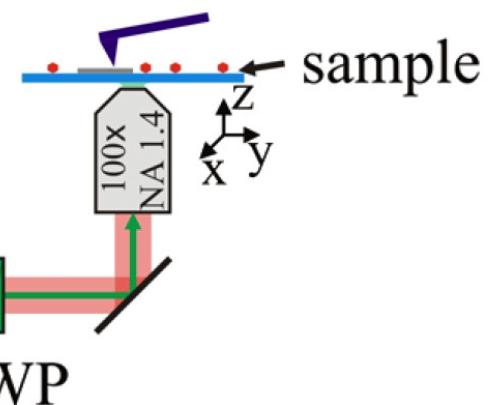
Chang et al., Phys. Rev. Lett. 97, 053002 (2006)

Previous experiments:
Akimov et al., Nature (London) 450, 402 (2007).
R. Kolesov et al., Nature Phys. 5, 470 (2009).
No full control

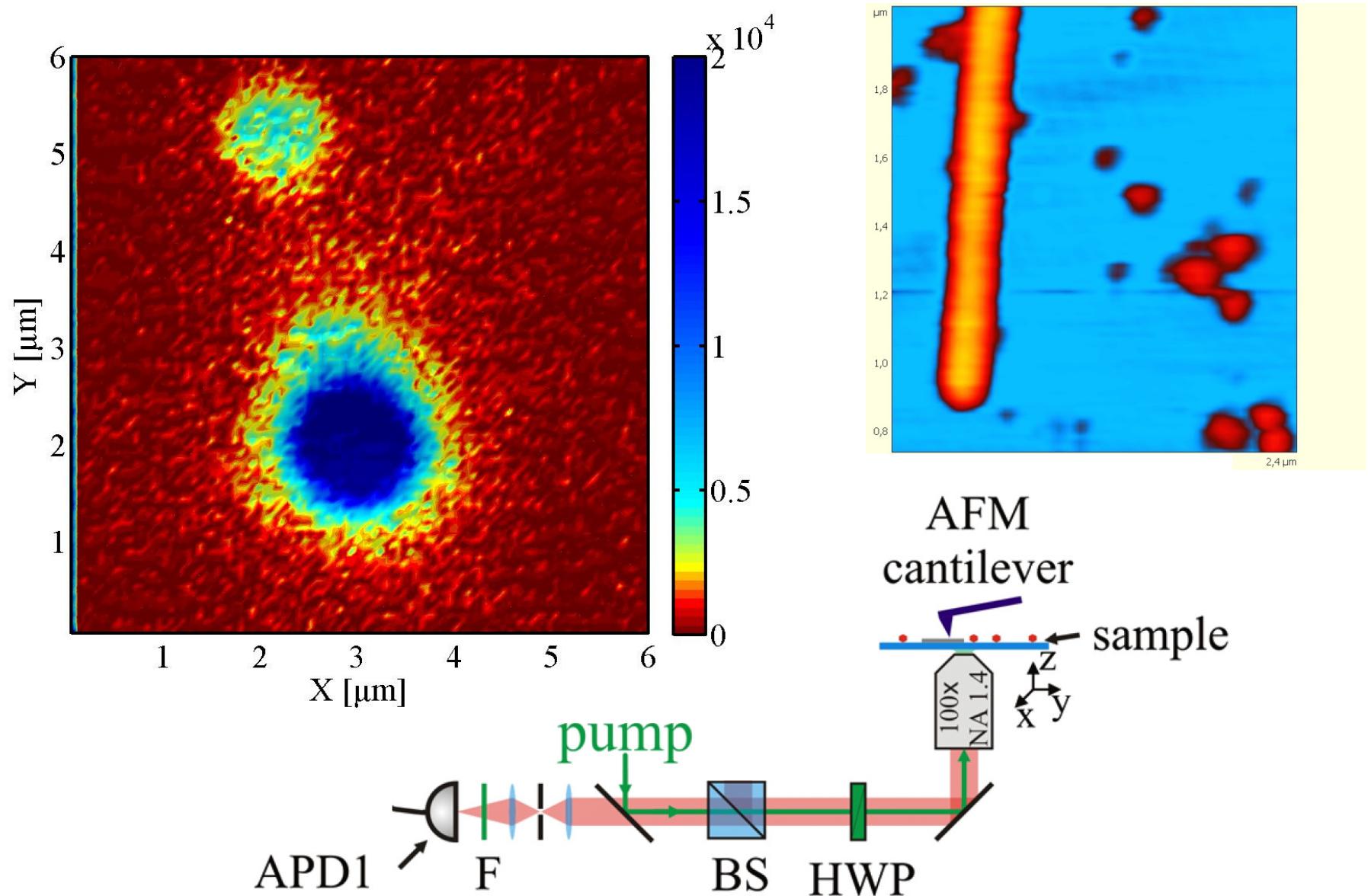
Integrated Confocal and AF microscope



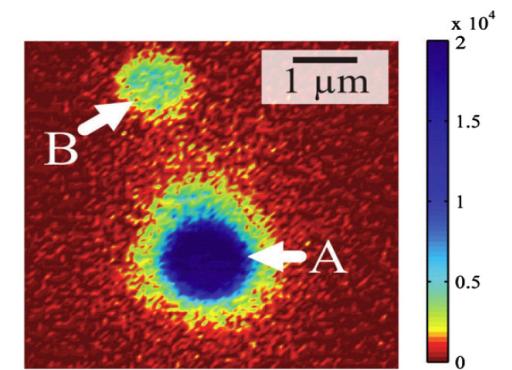
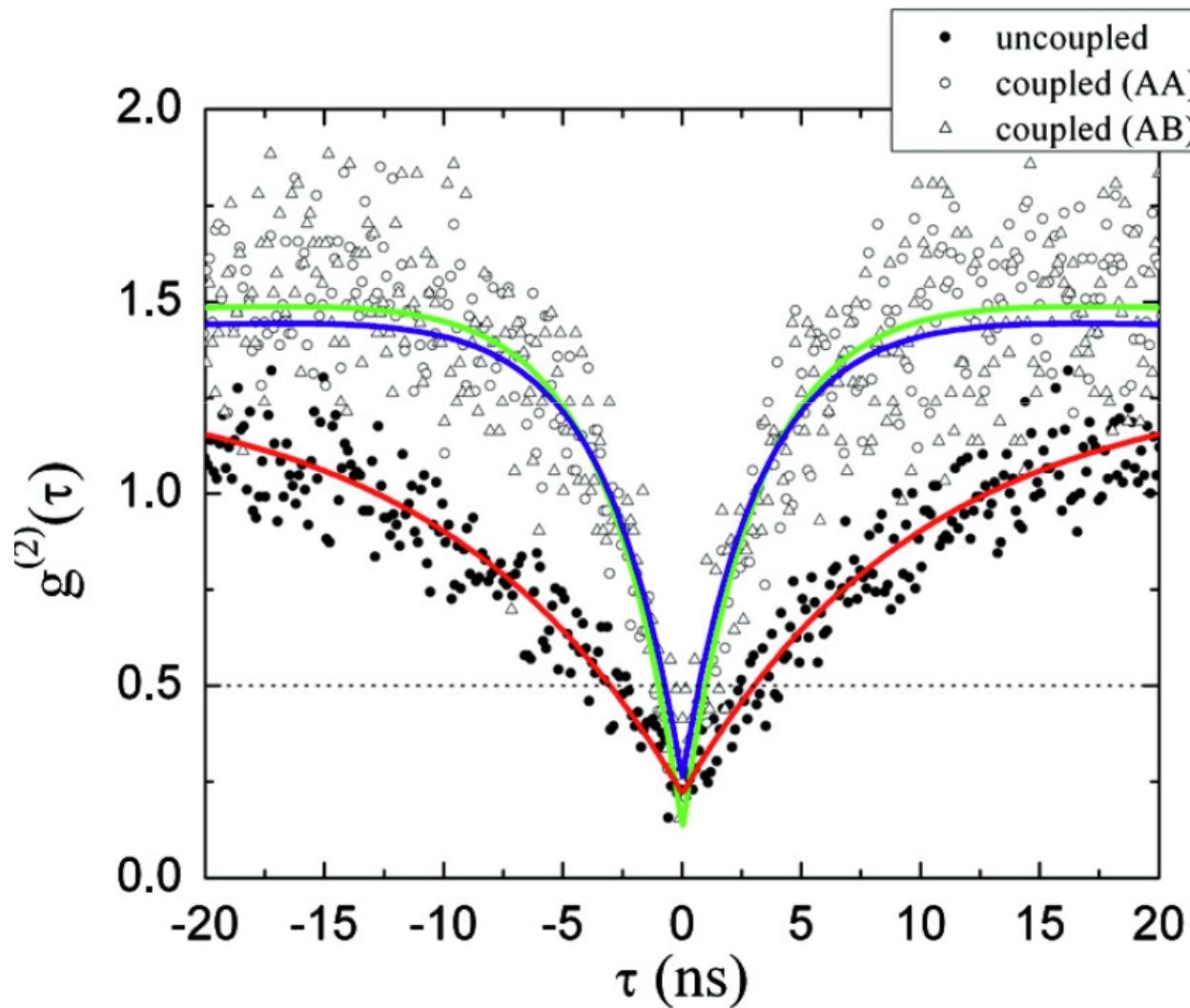
AFM
cantilever



Increasing the decay rate by a factor of 3.6



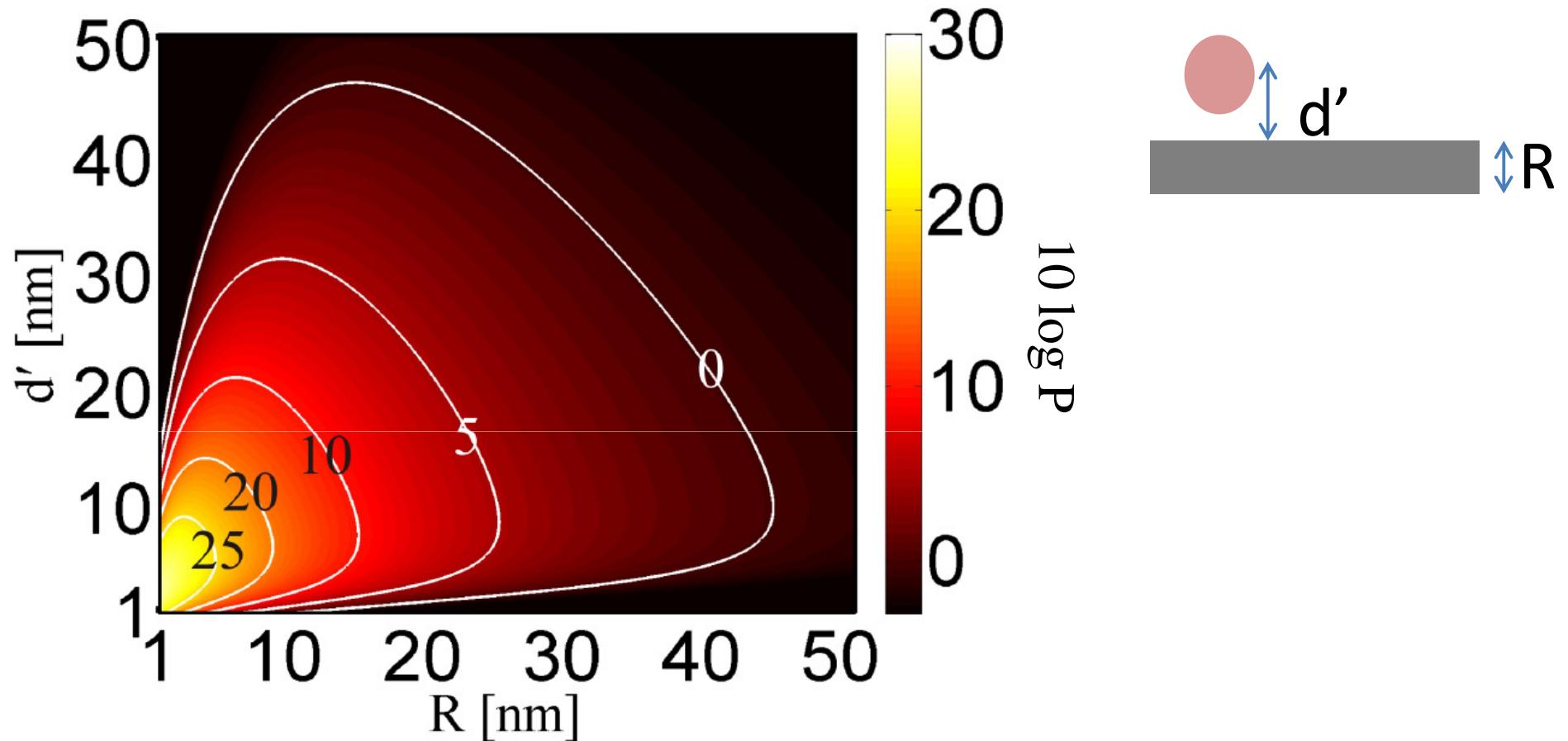
Single plasmons are generated



Single plasmon propagation:

$$g^{(2)}(0) < 0.5$$

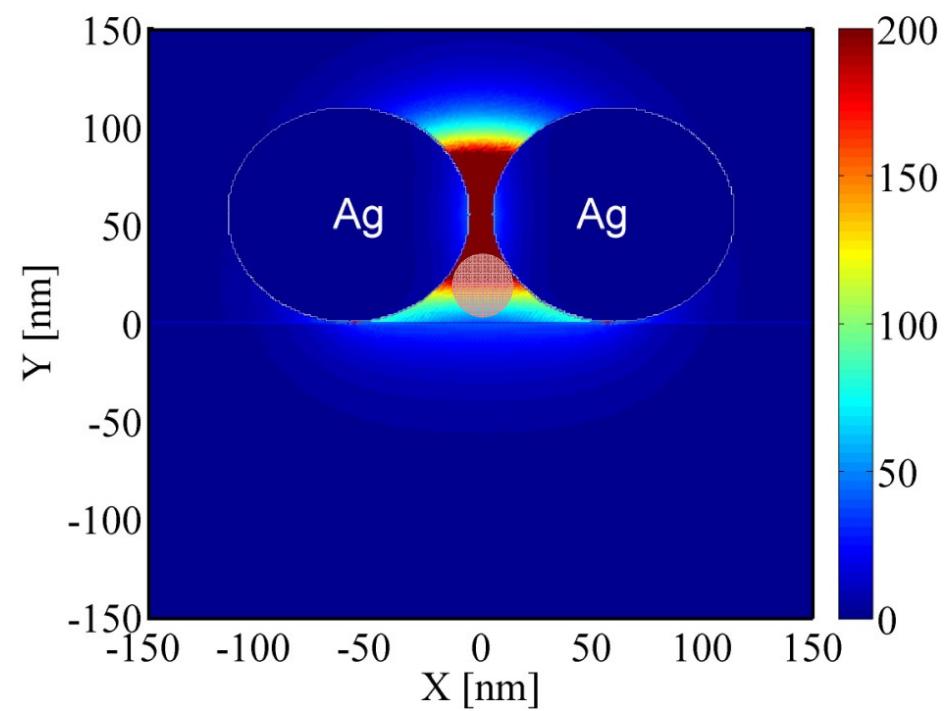
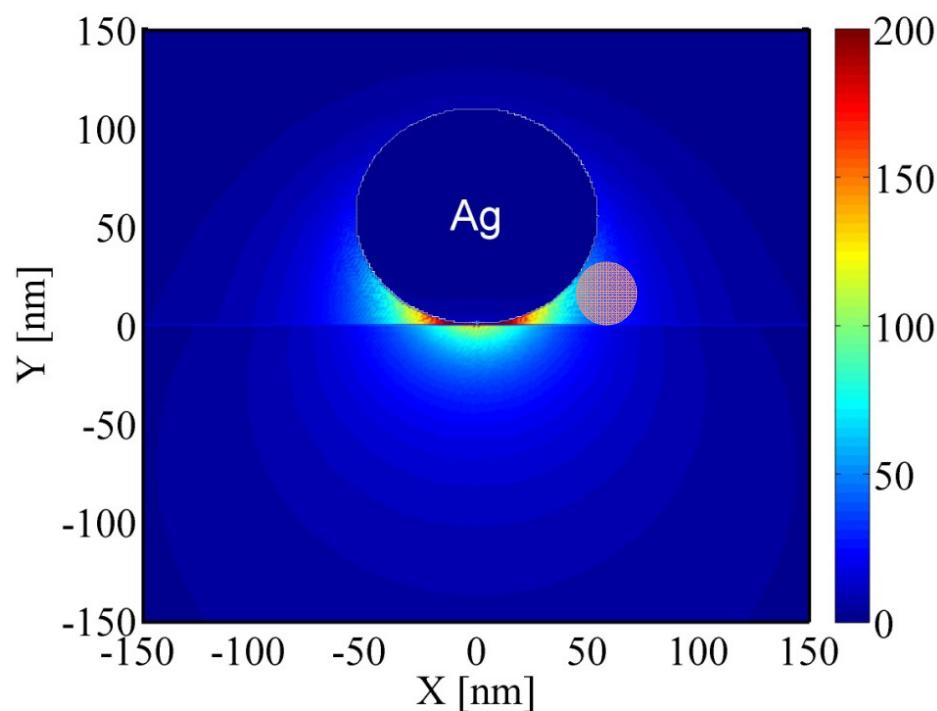
How do we get stronger coupling?

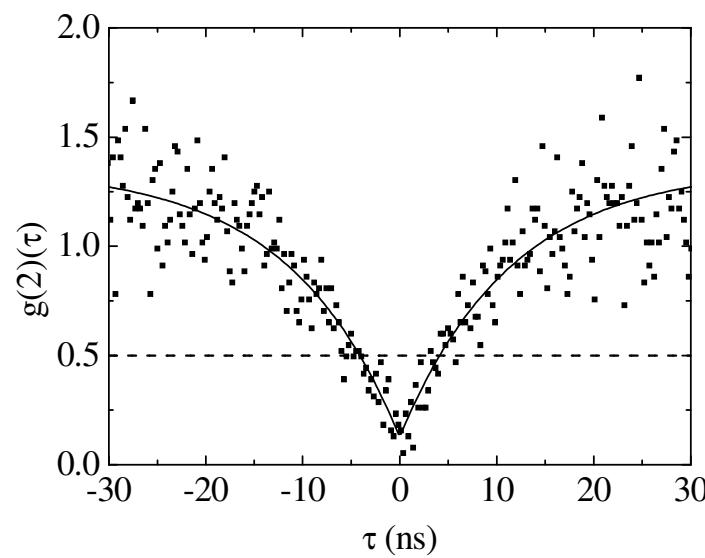
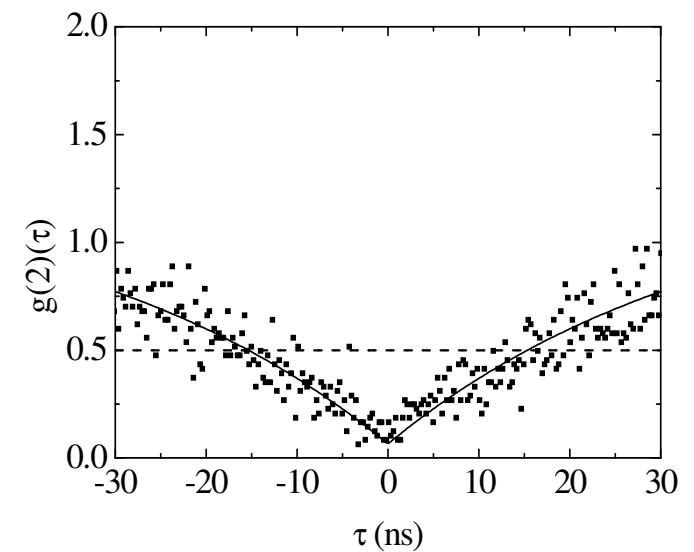
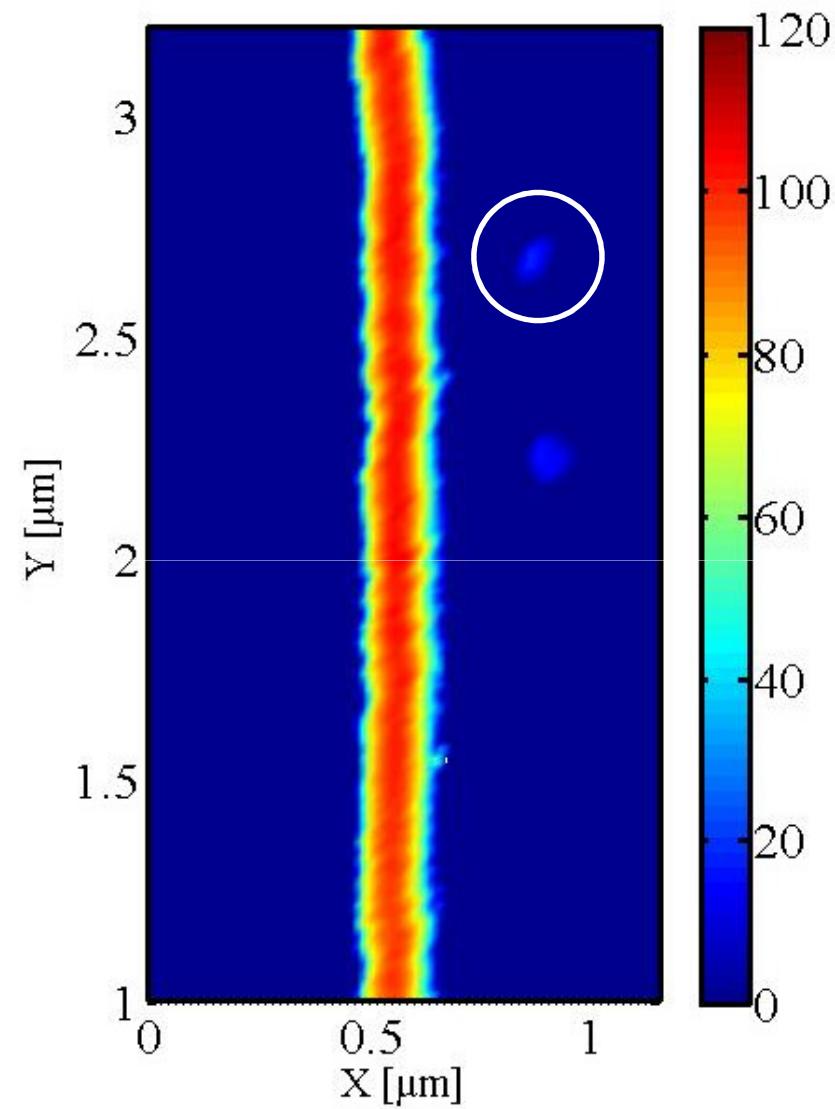


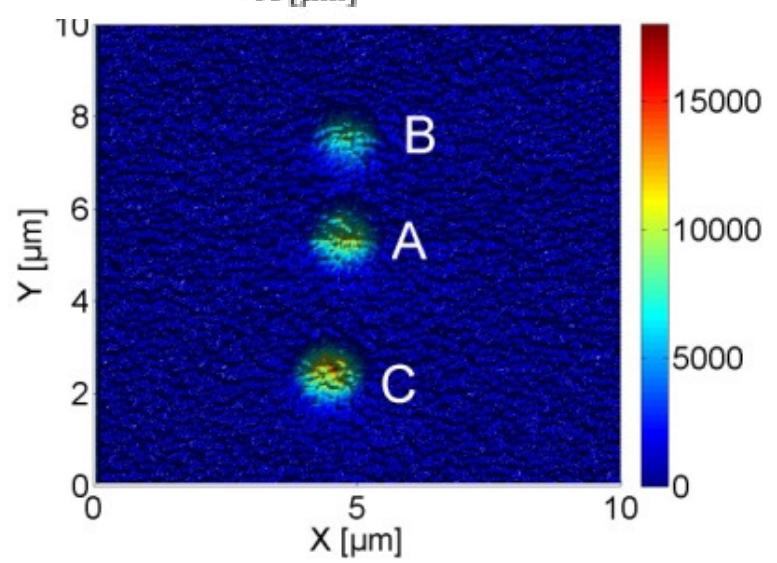
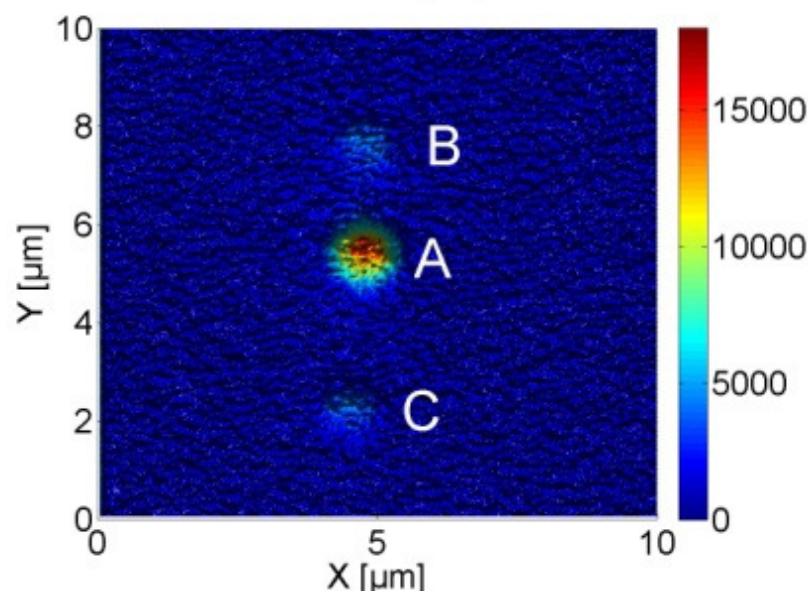
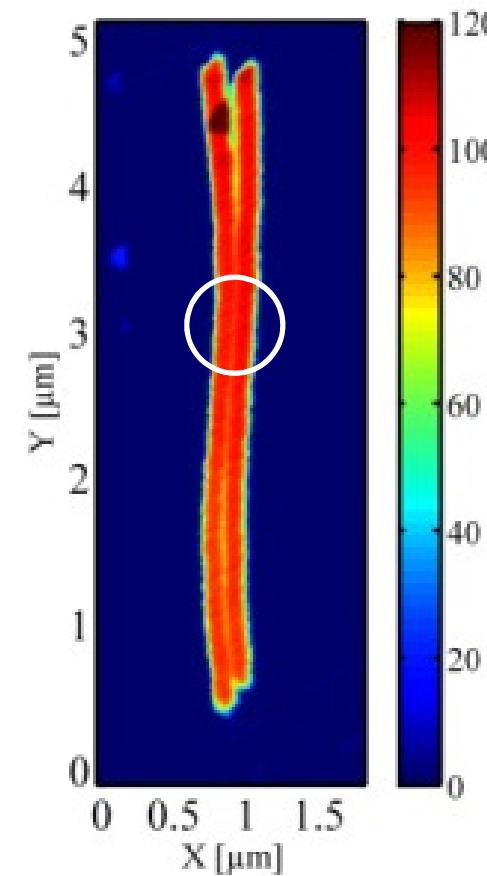
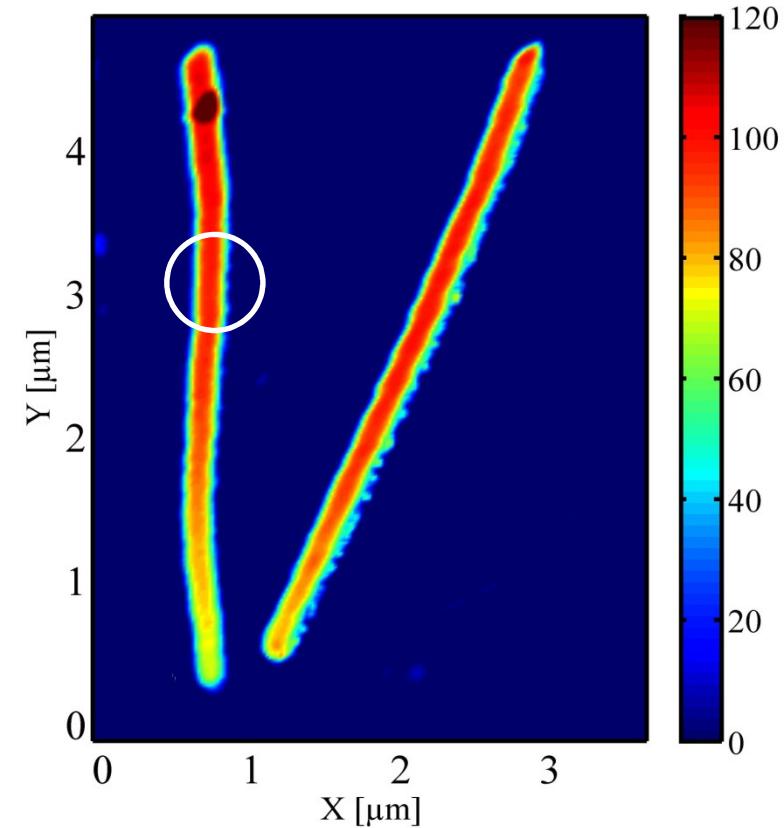
Smaller wires but not too small!

Smaller diamonds but not too small!

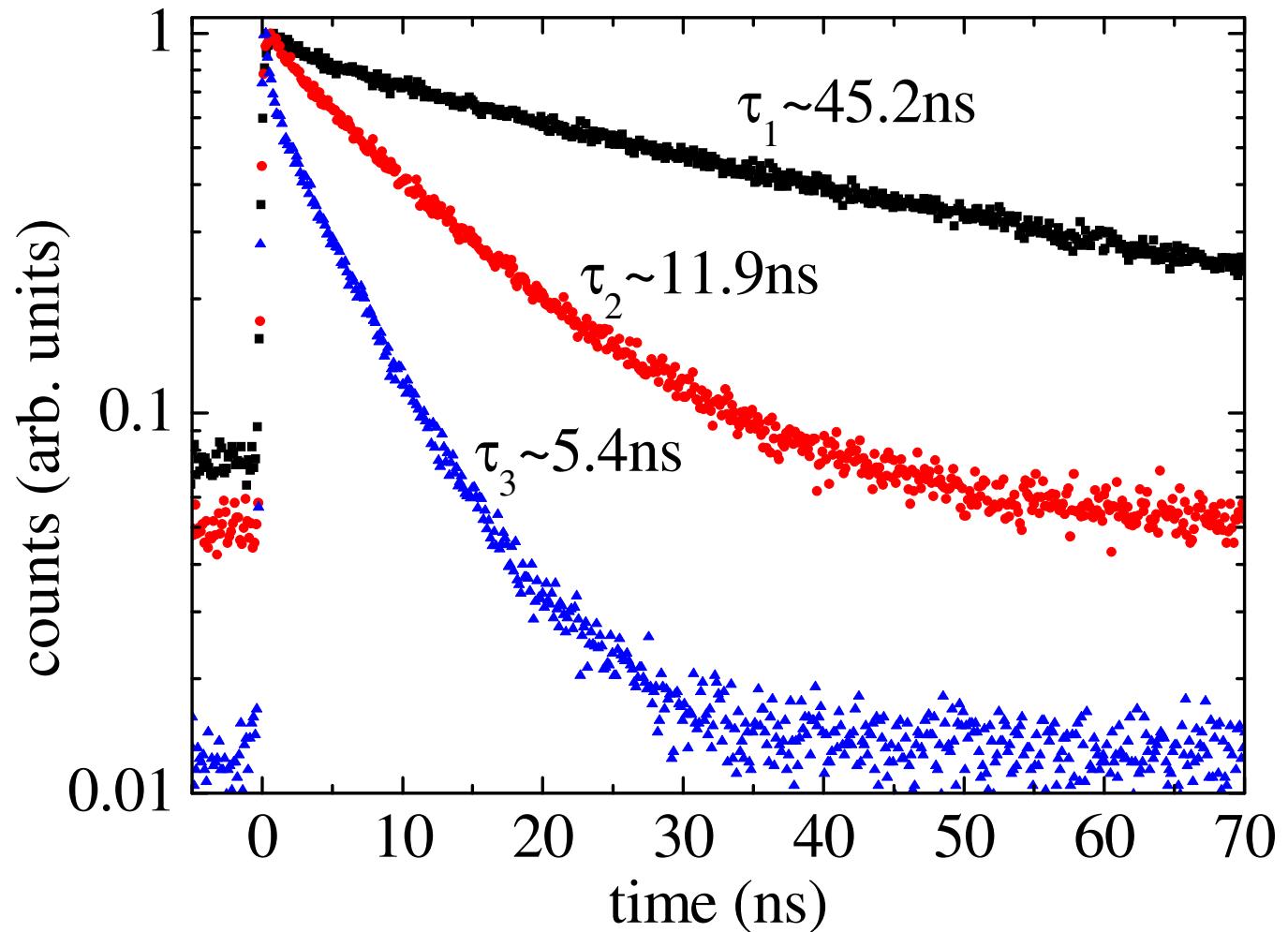
How to access the strong field

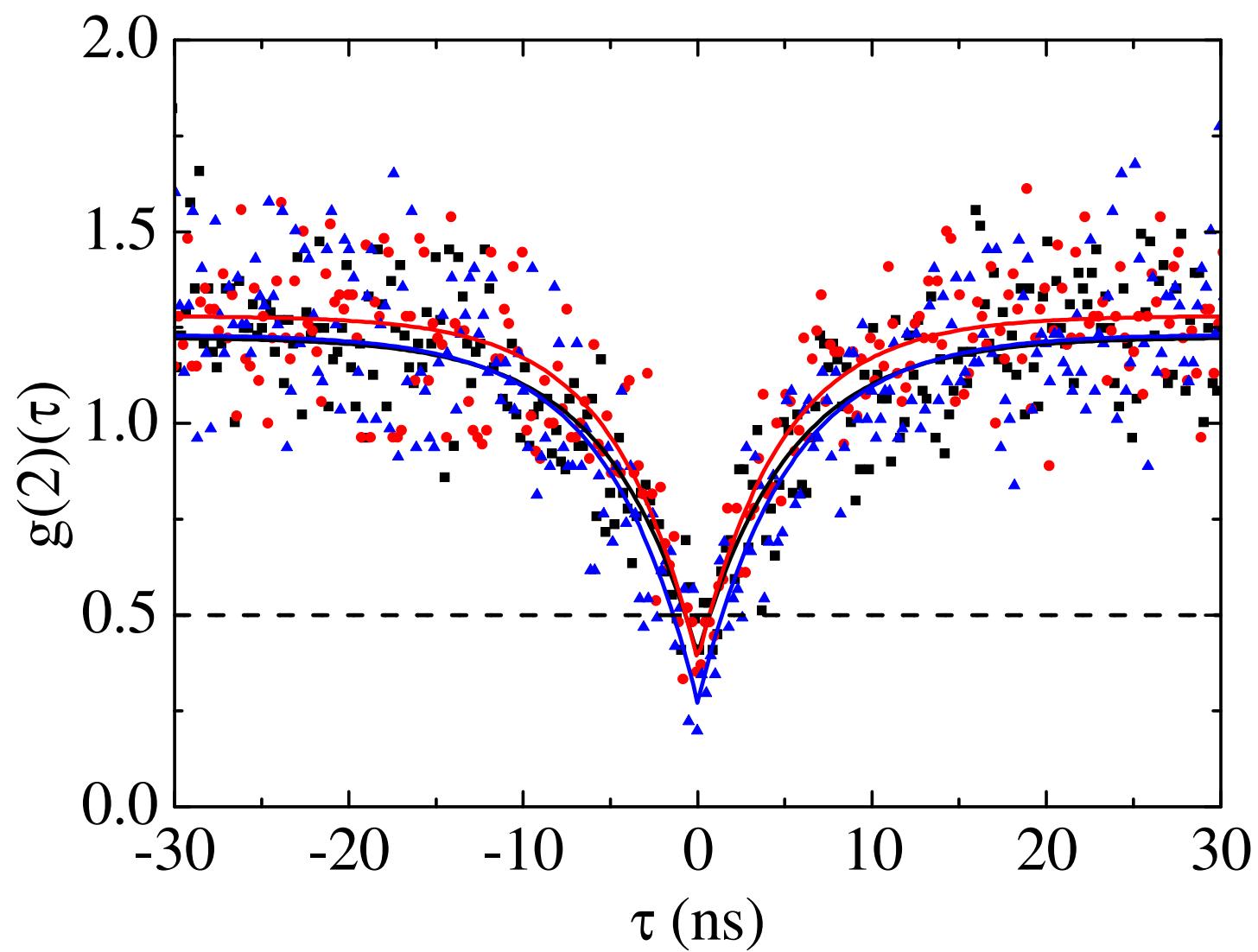






Increasing the decay rate by a factor of 8.3

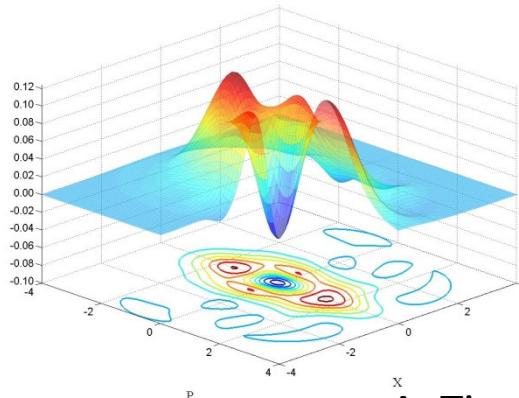




Towards large non-linearity

Measurement induced nonlinearity: Hadamard gate

$|-\alpha\rangle \rightarrow$



A. Tipsmark et al [arXiv:1107.0822](https://arxiv.org/abs/1107.0822)

Strong coupling



Max rate enhancement of 8.3



Huck et al, PRL 106, 096801 (2011)

What next?

- Tailored structures
- Smaller diamonds
- Resonant coupling at low temperature